

## CHANDRA STUDY OF AN OVERDENSITY OF X-RAY SOURCES AROUND TWO DISTANT ( $z \sim 0.5$ ) CLUSTERS

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### ABSTRACT

We present results from a *Chandra X-Ray Observatory* study of the field X-ray source populations in four different observations: two high-redshift ( $z \sim 0.5$ ) clusters of galaxies 3C 295 and RX J003033.2+261819; and two noncluster fields with similar exposure time. Surprisingly, the 0.5–2 keV source surface densities ( $\sim 900$ – $1200$  sources  $\text{deg}^{-2}$  at a flux limit of  $1.5 \times 10^{-15}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$ ) measured in an  $\sim 8' \times 8'$  area surrounding each cluster exceed by a factor of  $\sim 2$  the value expected on the basis of the *ROSAT* and *Chandra* log  $N$ -log  $S$ , with a significance of  $\sim 2 \sigma$  each, or  $\sim 3.5 \sigma$  when the two fields are combined (i.e., a probability to be a statistical fluctuation of  $< 1\%$  and  $< 0.04\%$ , respectively). The same analysis performed on the noncluster fields and on the outer chips of the cluster fields does *not* show evidence of such an excess. In both cluster fields, the summed 0.5–10 keV spectrum of the detected objects is well fitted by a power law with  $\Gamma \sim 1.7$  similar to active galactic nuclei (AGNs) and shows no sign of intrinsic absorption. The few ( $\sim 10$  of 35) optical identifications available to date confirm that most of them are, as expected, AGNs, but the number of redshifts available is too small to allow conclusions on their nature. We discuss possible interpretations of the overdensity in terms of a statistical variation of cosmic background sources; a concentration of AGNs and/or powerful starburst galaxies associated with the clusters; and gravitational lensing of background QSOs by the galaxy clusters. All explanations, however, are difficult to reconcile with the large number of excess sources detected. Deeper X-ray observations and more redshifts measurements are clearly required to settle the issue.

*Subject headings:* galaxies: active — galaxies: clusters: general — X-rays: galaxies — X-rays: general

### 1. INTRODUCTION

Since its launch date on 1999 July 23, the *Chandra X-Ray Observatory* has performed a number of pointed observations aimed at verifying the satellite functioning and at calibrating the instrument responses. This paper reports the analysis of serendipitous sources detected in three of these observations.

Among the most remarkable characteristics of *Chandra* are its unprecedented sensitivity and spatial resolution ( $\lesssim 1''$ ) over the entire 0.1–10 keV band (Van Speybroeck et al. 1997), a factor of  $\sim 10$  better than any previous X-ray mission. This provides an order-of-magnitude advance in detecting faint point sources ( $\sim 10$ – $100$  times fainter than *ROSAT* and *ASCA* at a given exposure time) because of the 100 times reduced background per beam element. On the basis of the *ROSAT* measurements (Hasinger et al. 1998), at a 0.5–2 keV flux limit of  $3 \times 10^{-15}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$ , the source density in *Chandra* observations is expected to be  $340 \pm 30$   $\text{deg}^{-2}$ , giving approximately six sources per chip. This opens new possibilities for detecting many serendip-

itous X-ray sources even in observations with modest ( $\sim$  few tens ks) exposures. This will lead to the collection of sufficiently large samples to enable detailed study of the log  $N$ -log  $S$ , the X-ray background, and potentially the spatial distribution of sources to map large-scale structure.

This paper reports on the serendipitous sources in *Chandra* observations of two medium- $z$  clusters RX J003033.2+261819 ( $z = 0.5$ , Vikhlinin et al. 1998, “RX J0030” here after) and 3C 295 ( $z = 0.46$ , Dressler & Gunn 1992), which suggest an excess number of serendipitous X-ray sources compared to a noncluster field and to the predictions based on the *ROSAT* and *Chandra* (0.5–2 keV) log  $N$ -log  $S$  measurements. The clusters have 0.5–10 keV luminosities of  $\sim 10^{44}$  and  $10^{45}$  ergs  $\text{s}^{-1}$  and temperatures of  $\sim 4$  and 4.4 keV, respectively (W. Forman et al., 2000, in preparation; Harris et al. 2000). We first show (§ 4.1) the source densities obtained from the on-axis chips of the RX J0030 and 3C 295 fields of view (FOVs) and compare them to densities in the outer chips and to two comparison fields (§§ 4.2 and 4.3): one obtained when *Chandra* pointed away from the radiant of the 1999 Leonid meteor shower (hereafter anti-Leonid) and one obtained from a calibration observation of 3C 273. The average X-ray spectral properties and available optical identifications of these sources are given in § 5. Possible interpretations of these results are discussed in § 6 and conclusions are reported in § 7.  $H_0 = 50$   $\text{km s}^{-1}$   $\text{Mpc}^{-1}$  and  $q_0 = 0.5$  are used throughout.

### 2. OBSERVATIONS AND DATA REDUCTIONS

The *Chandra X-Ray Observatory* (Weisskopf, O’Dell, & Van Speybroeck 1996) consists of four pairs of concentric

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TABLE 1  
LOG OF OBSERVATIONS

FOV	OBSERVATION IDENTIFICATION	OBSERVATIONAL MODE	EXPOSURE <sup>a</sup> (s)	NOMINAL POINTING	
				R.A. (2000)	Decl. (2000)
Cluster Fields					
RX J0030 .....	1190	ACIS-S, 4 chips	15375	00 30 40	+26 18 00
	1226	ACIS-S, 4 chips	14746	...	...
	merged	ACIS-S, 4 chips	30121	...	...
3C 295 .....	578	ACIS-S, 4 chips	18280	14 11 10	+52 13 01
Comparison Fields					
Anti-Leonid.....	1479	ACIS-I, 4 chips	20580	22 13 12	-22 10 41
3C 273 .....	1712	ACIS-S, 1 chip (S3)	22800	12 29 06	+02 03 14

NOTE.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.

<sup>a</sup> Exposure time obtained for the on-axis chips (ACIS-S3 for RX J0030, 3C 295, and 3C 273; and ACIS-I3 for anti-Leonid).

Wolter I mirrors reflecting 0.1–10 keV X-rays (Van Speybroeck et al. 1997) into one of the four focal plane detectors (ACIS-I/S or HRC-I/S). All the data presented in the following were taken from the *Chandra* public archive (G. Fabbiano et al., 2000, in preparation; see also <http://asc.harvard.edu/cda/>). The observations of RX J0030 and 3C 295 were performed with the ACIS-S, with the clusters lying within a few arcseconds of the optical axis location on the back-illuminated (BI) S3 chip. Chips S1–S4, I2 and I3 constitute the entire activated field of view. The anti-Leonid observation was performed with the ACIS-I configuration (i.e., with the focus nearly at the center of four front-illuminated (FI) CCDs—see the “*Chandra* Proposer’s Observatory Guide” 1999).<sup>8</sup> The observation of 3C 273 was performed with an ACIS-S(1–6) configuration. These two observations were chosen from the public archival data as the best available comparison fields because of their long exposures and high Galactic latitudes.

Details of the cleaning and reduction of the data are given in Appendix A. In total, we obtained “good” 0.1–10 keV data from four chips for each FOV, each chip having dimensions of 8′ × 8′. A log of the observations is given in Table 1.

### 3. ANALYSIS: THE SOURCE DETECTIONS, COUNTS, AND FLUXES

#### 3.1. The Source Detections

To localize the serendipitous source candidates in the fields, we applied a source detection algorithm in the *Chandra* Interactive Analysis of Observations (CIAO) (M. Elvis et al., 2000, in preparation) software: *wavdetect* (Freeman et al. 2000; Dobrzycki et al. 1999).<sup>9</sup> Source detection is easy with *Chandra* because the background is very uniform and low, even in the vicinity of the clusters, and one can therefore “see” the sources unambiguously. With the conservative threshold applied here (see Appendix B), all sources are indeed visible by eye. Details of the algorithm

and the procedure applied for the detections are given in Appendix B.

In the 0.5–2 keV energy band, a total of 53 and 44 point sources/four-chip FOV (clusters excluded) are found in the RX J0030 and 3C 295 fields, respectively. All detected sources were consistent with point sources. More sources are expected in the RX J0030 field because we have about 50% more usable exposure time. In the hard 2–10 keV energy band, these numbers are reduced to 13 and five, respectively. If we restrict ourselves to the central S3 chip (where systematic effects are expected to be smaller), we find 23 and 17 sources between 0.5–2 keV, and six and four sources between 2 and 10 keV, respectively. These sources are listed in Tables 2 and 3. For comparison, we also list the sources detected in the whole four chips of the anti-Leonid field and in the ACIS-S3 chip of the 3C 273 field (Tables 4 and 5).

Figure 1 shows an overlay of the detected sources with the X-ray images between 0.5 and 2 keV, where the images have been smoothed using a Gaussian function with  $\sigma = 1$  pixel = 2″. All the detected sources are clearly visible in the images.

#### 3.2. The Source Counts and Fluxes

Source count rates were obtained using the *wavdetect* algorithm from regions with typical radii of  $\sim 3''$  (on-axis) and  $10''$  (off-axis). The measured counts were first corrected for vignetting and then converted to an emitted, unabsorbed flux. A description of this procedure is given in Appendix C.

Tables 2 and 3 report the 0.5–2 keV and 2–10 keV measured fluxes for the detected sources (the sixth column in each table) and, for sources with known redshifts, the corresponding luminosities. The total 0.5–2 keV flux of the point sources in RX J0030 (most of which are within a 5′ radius) is  $2.2 \times 10^{-13}$  ergs cm<sup>-2</sup> s<sup>-1</sup>,  $\sim 1.6$  times larger than the cluster flux over the same energy band. About half of this flux is in the  $z = 0.492$  CRSS QSO (source 1 in Table 2). In 3C 295, the point sources sum to  $1.6 \times 10^{-13}$  ergs cm<sup>-2</sup> s<sup>-1</sup>, about one-third of the cluster flux.

## 4. DENSITIES OF THE SERENDIPITOUS SOURCES

### 4.1. In the Central Chips (S3)

Here we focus on the results obtained from the central

<sup>8</sup> “*Chandra* Proposer’s Observatory Guide,” 1999, version 2.0, 2000 March 15, <http://asc.harvard.edu/udocs/docs/docs.html>.

<sup>9</sup> A. Dobrzycki, H. Ebeling, K. Glotfelty, P. Freeman, F. Damiani, M. Elvis, & T. Calderwood, 1999, “*Chandra* DETECT 1.0 User Guide,” rev. 1.0, <http://asc.harvard.edu/udocs/docs/docs.html>.

TABLE 2  
LIST OF POINT SOURCES DETECTED BETWEEN 0.5 AND 2 keV AND 2 keV AND 10 keV IN THE RX J0030 FIELD—ACIS-S3 CENTRAL CHIP (TARGET CLUSTER EXCLUDED)

Row	R.A. (2000)	Decl. (2000)	Net Counts	Background Counts	Flux ( $10^{-15}$ cgs)	$\theta_{\text{off}}$ (arcmin)	$B$	$\alpha_{\text{ox}}^a$	$z$	$L_X^b$ ( $10^{42}$ ergs $\text{s}^{-1}$ )	$M_B$	Class	Reference
0.5–2 keV													
1	00 30 39.5	+26 20 56.4	1023.49 ± 32.47	6.76 ± 0.03	112.58 ± 3.57	3.0	17.3	1.62	0.492	134.2	-25.2	QSO	1
2	00 30 57.9	+26 17 44.8	120.54 ± 11.49	4.01 ± 0.02	13.26 ± 1.26	5.0	...	...	...	...	...	...	...
3	00 30 47.9	+26 16 48.6	96.23 ± 10.17	4.85 ± 0.02	10.59 ± 1.12	3.0	18.4	1.85	0.246	2.9	-22.6	Sey2	2
4	00 30 51.4	+26 17 14.3	84.14 ± 9.63	5.54 ± 0.03	9.26 ± 1.06	3.6	19.7	1.66	0.129	0.7	-19.6	Galaxy?	3
5	00 30 26.1	+26 16 49.7	83.06 ± 9.36	3.20 ± 0.02	9.14 ± 1.03	2.6	19.2	1.76	0.269	3.0	-21.9	QSO	3
6	00 30 27.8	+26 13 60.0	49.97 ± 7.50	3.41 ± 0.02	5.50 ± 0.82	4.6	19.8	1.79	...	...	...	...	...
7	00 30 27.9	+26 15 15.2	38.33 ± 6.67	4.77 ± 0.02	4.22 ± 0.73	3.4	...	...	...	...	...	...	...
8	00 30 41.7	+26 17 42.7	38.04 ± 6.44	3.69 ± 0.02	4.18 ± 0.71	1.3	...	...	...	...	...	...	...
9	00 30 33.3	+26 14 52.4	37.59 ± 6.50	3.42 ± 0.02	4.14 ± 0.71	3.3	19.3	1.86	0.247	1.1	-21.3	Galaxy	3
10	00 30 44.4	+26 18 32.6	37.45 ± 6.43	3.71 ± 0.02	4.12 ± 0.71	2.0	...	...	...	...	...	...	...
11	00 30 31.0	+26 16 00.6	35.43 ± 6.24	2.98 ± 0.02	3.90 ± 0.69	2.4	...	...	...	...	...	...	...
12	00 30 59.8	+26 17 18.9	34.80 ± 6.30	2.34 ± 0.02	3.83 ± 0.69	5.5	...	...	...	...	...	...	...
13	00 30 54.8	+26 14 47.5	28.81 ± 6.02	4.94 ± 0.02	3.17 ± 0.66	5.4	...	...	...	...	...	...	...
14	00 30 37.2	+26 19 04.4	25.79 ± 5.68	6.91 ± 0.03	2.84 ± 0.62	1.1	...	...	...	...	...	...	...
15	00 30 51.8	+26 18 56.6	25.34 ± 5.47	3.51 ± 0.02	2.79 ± 0.60	3.7	...	...	...	...	...	...	...
16	00 30 21.7	+26 18 34.4	23.93 ± 5.24	2.72 ± 0.02	2.63 ± 0.58	3.3	...	...	...	...	...	...	...
17	00 30 44.3	+26 17 02.7	23.62 ± 5.23	3.56 ± 0.02	2.60 ± 0.58	2.1	20.2	1.81	...	...	...	...	...
18	00 30 49.6	+26 16 05.1	19.93 ± 4.96	3.73 ± 0.02	2.19 ± 0.55	3.7	...	...	...	...	...	...	...
19	00 30 42.7	+26 17 48.6	19.60 ± 4.68	2.33 ± 0.02	2.16 ± 0.51	1.5	...	...	...	...	...	...	...
20	00 30 39.7	+26 14 19.5	18.15 ± 4.64	2.51 ± 0.02	2.00 ± 0.51	3.9	19.1	2.03	...	...	...	...	...
21	00 30 27.8	+26 15 28.9	14.61 ± 4.23	2.76 ± 0.02	1.61 ± 0.47	3.2	...	...	...	...	...	...	...
22	00 30 34.6	+26 13 51.9	13.46 ± 4.18	3.11 ± 0.02	1.48 ± 0.46	4.3	...	...	...	...	...	...	...
23	00 30 31.2	+26 14 38.3	11.84 ± 3.87	2.54 ± 0.02	1.30 ± 0.43	3.6	...	...	...	...	...	...	...
2–10 keV													
1	00 30 39.5	+26 20 56.2	596.50 ± 25.10	19.01 ± 0.05	417.55 ± 17.57	...	...	...	...	497.7	...	...	...
4	00 30 51.4	+26 17 14.8	37.38 ± 7.23	12.81 ± 0.04	26.16 ± 5.06	...	...	...	...	34.4	...	...	...
3	00 30 47.9	+26 16 48.9	32.51 ± 6.86	13.19 ± 0.04	22.76 ± 4.80	...	...	...	...	6.3	...	...	...
11	00 30 31.0	+26 16 00.5	28.10 ± 6.40	12.23 ± 0.04	19.67 ± 4.48	...	...	...	...	...	...	...	...
2	00 30 57.8	+26 17 44.8	16.44 ± 4.96	6.45 ± 0.03	11.51 ± 3.47	...	...	...	...	...	...	...	...
5	00 30 26.2	+26 16 49.6	14.60 ± 4.65	6.61 ± 0.03	10.22 ± 3.26	...	...	...	...	13.5	...	...	...

NOTE.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds. Position uncertainties are estimated to be within  $\sim 2''$ .  
<sup>a</sup>  $\alpha_{\text{ox}} = -\log(S_{2\text{keV}}/S_{2.500\text{\AA}})/2.605$ , as defined in Stocke et al. 1991.  
<sup>b</sup>  $z = 0.5$  and  $z = 0.46$  are assumed for RX J0030 and 3C 295, respectively, unless the redshift is available from the literature.  
REFERENCES.—(1) Ciliegi et al. 1995; (2) Boyle et al. 1995; (3) Brandt et al. 2000.

TABLE 3  
LIST OF POINT SOURCES DETECTED BETWEEN 0.5 AND 2 keV AND 2 AND 10 keV IN THE 3C 295 FIELD—ACIS-S3 CENTRAL CHIP (TARGET CLUSTER EXCLUDED)

Row	R.A. (2000)	Decl. (2000)	Net Counts	Background Counts	Flux ( $10^{-15}$ ergs)	$\theta_{\text{off}}$ (arcmin)	$B$	$\alpha_{\text{ox}}^a$	$z$	$L_x^b$ ( $10^{42}$ ergs $s^{-1}$ )	$M_B$	Class	Reference
0.5–2 keV													
1	14 11 23.4	+ 52 13 31.4	360.33 ± 19.04	3.50 ± 0.02	55.85 ± 2.95	1.5	19.8	1.34	0.47	58.6	-22.6	Sey1	1
2	14 11 04.1	+ 52 17 56.0	121.50 ± 11.59	2.56 ± 0.02	18.83 ± 1.80	6.3	18.6	1.73	...	...	-23.6	...	...
3	14 11 19.5	+ 52 14 00.0	112.82 ± 10.85	4.50 ± 0.02	17.49 ± 1.68	1.9	22.0	1.20	1.29	172.0	-22.8	QSO	2
4	14 11 24.8	+ 52 19 23.9	82.36 ± 9.65	2.11 ± 0.02	12.77 ± 1.50	7.4	...	...	...	...	...	...	...
5	14 11 27.4	+ 52 11 31.2	51.48 ± 7.38	3.25 ± 0.02	7.98 ± 1.14	1.4	20.1	1.62	0.59	13.9	-22.8	Sbc	3
6	14 10 59.5	+ 52 11 52.8	51.15 ± 7.31	1.06 ± 0.01	7.93 ± 1.13	3.1	...	...	...	...	...	...	...
7	14 11 30.8	+ 52 14 23.4	35.40 ± 6.12	1.31 ± 0.01	5.49 ± 0.95	2.9	13.7	2.68	...	...	-28.5	Star?	...
8	14 11 13.7	+ 52 13 40.4	33.12 ± 5.92	1.92 ± 0.02	5.13 ± 0.92	1.8	21.2	1.53	2.08	147.0	-24.8	QSO	1
9	14 11 09.0	+ 52 16 45.1	30.34 ± 5.79	1.28 ± 0.01	4.70 ± 0.90	4.9	...	...	...	...	...	...	...
10	14 10 57.4	+ 52 11 30.0	24.73 ± 5.16	1.01 ± 0.01	3.83 ± 0.80	3.5	...	...	...	...	...	...	...
11	14 11 14.5	+ 52 16 10.9	22.44 ± 5.00	1.48 ± 0.01	3.48 ± 0.78	4.2	...	...	...	...	...	...	...
12	14 11 20.7	+ 52 14 11.5	22.06 ± 4.93	2.08 ± 0.02	3.42 ± 0.76	2.1	...	...	...	...	...	...	...
13	14 11 29.2	+ 52 13 32.5	20.11 ± 4.72	2.02 ± 0.02	3.12 ± 0.73	2.1	21.5	1.56	0.66	7.0	-21.7	Scd	3
14	14 11 26.3	+ 52 18 47.8	16.33 ± 4.36	1.01 ± 0.01	2.53 ± 0.68	6.8	...	...	...	...	...	...	...
15	14 11 19.8	+ 52 15 46.3	12.28 ± 3.73	1.13 ± 0.01	1.90 ± 0.58	3.7	...	...	...	...	...	...	...
16	14 11 33.1	+ 52 11 02.6	10.93 ± 3.51	1.20 ± 0.01	1.69 ± 0.54	2.4	...	...	...	...	...	...	...
17	14 11 27.0	+ 52 14 14.0	8.24 ± 3.20	1.87 ± 0.01	1.28 ± 0.50	2.4	...	...	...	...	...	...	...
2–10 keV													
1	14 11 23.4	+ 52 13 31.2	80.04 ± 9.14	3.69 ± 0.02	91.2 ± 10.4	...	...	...	...	98.5	...	...	...
2	14 11 04.1	+ 52 17 55.8	45.95 ± 7.87	10.49 ± 0.04	52.4 ± 9	...	...	...	...	...	...	...	...
5	14 11 27.4	+ 52 11 31.3	30.31 ± 5.80	3.53 ± 0.02	34.6 ± 6.6	...	...	...	...	60.6	...	...	...
3	14 11 19.4	+ 52 14 00.4	29.40 ± 5.93	5.69 ± 0.03	33.5 ± 6.8	...	...	...	...	329.1	...	...	...

NOTE.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds. Position uncertainties are estimated to be within  $\sim 2''$ .  
 $\alpha_{\text{ox}}^a = -\log(S_{2\text{ keV}}/S_{2500\text{ \AA}})/2.605$ , as defined in Stocke et al. 1991.  
 $z = 0.5$  and  $z = 0.46$  are assumed for RX J0030 and 3C 295, respectively, unless the redshift is available from the literature.  
REFERENCES.—(1) Dressler & Gunn 1992; (2) Hewitt & Burbidge 1989; (3) Thimm et al. 1994.

TABLE 4  
LIST OF POINT SOURCES DETECTED BETWEEN 0.5 AND 2 keV AND 2 AND 10 keV IN THE  
ANTI-LEONID FIELD—WHOLE FOV (4 CHIPS)

Row	R.A. (2000)	Decl. (2000)	Net Counts	Background Counts	Flux ( $10^{-15}$ cgs)	$\theta_{\text{off}}$ (arcmin)
0.5–2 keV						
1	22 13 52.4	–22 15 56.2	333.20 ± 19.86	5.38 ± 0.03	82.97 ± 4.95	11.0
2	22 13 23.2	–22 07 24.5	257.09 ± 16.37	1.20 ± 0.01	64.02 ± 4.08	6.2
3	22 12 57.5	–22 21 35.5	213.74 ± 16.10	4.92 ± 0.02	53.22 ± 4.01	4.8
4	22 13 13.1	–22 04 25.5	109.38 ± 10.90	1.18 ± 0.01	27.23 ± 2.71	4.5
5	22 12 49.1	–22 11 32.2	74.91 ± 9.01	0.93 ± 0.01	18.65 ± 2.24	4.1
6	22 12 55.8	–22 10 04.6	72.03 ± 8.69	0.67 ± 0.01	17.94 ± 2.16	5.5
7	22 12 40.1	–22 07 48.1	51.28 ± 7.61	0.83 ± 0.01	12.77 ± 1.90	3.8
8	22 13 26.2	–22 05 50.5	49.22 ± 7.39	1.87 ± 0.01	12.26 ± 1.84	2.3
9	22 12 25.1	–22 07 28.5	43.80 ± 7.50	3.62 ± 0.02	10.91 ± 1.87	2.4
10	22 13 28.7	–22 11 51.6	34.99 ± 6.08	0.43 ± 0.01	8.71 ± 1.51	1.4
11	22 13 19.5	–22 08 36.0	31.99 ± 5.72	0.37 ± 0.01	7.97 ± 1.42	5.9
12	22 13 33.1	–22 10 03.8	27.81 ± 5.48	0.64 ± 0.01	6.92 ± 1.36	1.4
13	22 12 43.4	–22 18 38.1	24.44 ± 5.45	0.95 ± 0.01	6.09 ± 1.36	3.1
14	22 12 51.5	–22 13 48.7	24.07 ± 5.16	0.65 ± 0.01	5.99 ± 1.29	5.3
15	22 12 54.2	–22 19 05.3	23.68 ± 5.49	2.30 ± 0.02	5.90 ± 1.37	4.6
16	22 13 25.9	–22 16 44.6	22.24 ± 5.03	0.73 ± 0.01	5.54 ± 1.25	3.2
17	22 13 37.9	–22 08 29.3	22.14 ± 4.95	0.49 ± 0.01	5.51 ± 1.23	4.8
18	22 12 58.1	–22 13 59.9	21.72 ± 4.85	0.49 ± 0.01	5.41 ± 1.21	3.1
19	22 12 47.7	–22 17 49.8	19.10 ± 4.96	2.23 ± 0.02	4.76 ± 1.24	3.3
20	22 13 27.6	–22 14 24.6	16.62 ± 4.26	0.42 ± 0.01	4.14 ± 1.06	5.1
21	22 13 41.0	–22 15 10.1	16.06 ± 4.32	0.61 ± 0.01	4.00 ± 1.08	1.2
22	22 13 04.5	–22 14 36.6	15.44 ± 4.08	0.35 ± 0.01	3.84 ± 1.02	3.9
23	22 13 00.5	–22 08 32.9	14.00 ± 3.85	0.39 ± 0.01	3.49 ± 0.96	4.8
24	22 13 30.2	–22 03 20.3	13.76 ± 4.01	0.64 ± 0.01	3.43 ± 1.00	6.1
25	22 13 12.6	–22 08 52.5	13.62 ± 3.72	0.32 ± 0.01	3.39 ± 0.93	5.6
26	22 13 29.0	–22 07 55.5	13.30 ± 3.78	0.33 ± 0.01	3.31 ± 0.94	8.0
27	22 12 54.5	–22 08 01.4	13.24 ± 3.80	0.44 ± 0.01	3.30 ± 0.95	8.6
28	22 13 03.7	–22 08 04.2	12.89 ± 3.69	0.40 ± 0.01	3.21 ± 0.92	8.0
29	22 12 57.5	–22 19 15.9	12.55 ± 4.12	2.00 ± 0.02	3.13 ± 1.02	8.3
30	22 13 20.6	–22 11 31.0	11.83 ± 3.50	0.28 ± 0.01	2.95 ± 0.87	8.2
31	22 13 07.9	–22 15 25.9	10.30 ± 3.36	0.32 ± 0.01	2.56 ± 0.84	5.5
32	22 12 38.1	–22 06 34.9	9.72 ± 3.56	1.37 ± 0.01	2.42 ± 0.89	7.3
33	22 13 08.1	–22 07 22.9	8.89 ± 3.07	0.33 ± 0.01	2.21 ± 0.77	7.2
34	22 13 06.5	–22 07 26.5	7.83 ± 2.90	0.37 ± 0.01	1.95 ± 0.72	9.9
35	22 13 49.1	–22 11 12.2	7.57 ± 2.97	0.25 ± 0.01	1.88 ± 0.74	9.6
36	22 13 45.6	–22 11 57.0	7.35 ± 2.94	0.38 ± 0.01	1.83 ± 0.73	1.1
37	22 13 20.3	–22 03 24.7	7.22 ± 2.90	0.42 ± 0.01	1.80 ± 0.72	1.1
38	22 13 06.9	–22 05 56.7	6.96 ± 2.78	0.37 ± 0.01	1.73 ± 0.69	3.3
39	22 13 20.8	–22 07 40.8	6.79 ± 2.72	0.39 ± 0.01	1.69 ± 0.68	9.8
40	22 13 18.4	–22 10 20.5	6.75 ± 2.63	0.21 ± 0.00	1.68 ± 0.66	8.9
2–10 keV						
2	22 13 23.2	–22 07 24.3	58.91 ± 7.97	2.22 ± 0.02	7.19 ± 0.97	...
3	22 12 57.4	–22 21 34.7	58.48 ± 9.44	14.46 ± 0.04	7.13 ± 1.15	...
4	22 13 13.0	–22 04 25.4	43.04 ± 7.04	2.93 ± 0.02	5.25 ± 0.86	...
40	22 13 18.4	–22 10 20.6	36.33 ± 6.13	1.46 ± 0.01	4.43 ± 0.75	...
5	22 12 49.1	–22 11 32.8	30.31 ± 5.95	2.65 ± 0.02	3.70 ± 0.73	...
12	22 13 33.1	–22 10 03.1	15.04 ± 4.35	2.75 ± 0.02	1.83 ± 0.53	...
23	22 12 54.6	–22 08 00.9	14.51 ± 4.08	1.23 ± 0.01	1.77 ± 0.50	...
6	22 12 57.4	–22 21 34.7	12.70 ± 3.89	1.78 ± 0.01	1.55 ± 0.47	...

NOTE.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds. Position uncertainties are estimated to be within  $\sim 2''$ .

chip (S3) of the RX J0030 and 3C 295 fields, which contain the clusters themselves, and compare them to the *ROSAT*, *ASCA* and *Chandra* log  $N$ -log  $S$ .

The faintest source in the RX J0030 field has a flux of  $\sim 1.3 \times 10^{-15}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$  (0.5–2 keV) and  $\sim 10^{-14}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$  (2–10 keV). For 3C 295, the faintest has

$\sim 1.3 \times 10^{-15}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$  and  $\sim 3.3 \times 10^{-14}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$  between 0.5 and 2 keV and 2 and 10 keV, respectively. All the sources detected at 2–10 keV were also detected in the 0.5–2 keV band.

We find that the minimum number of photon counts of the detected sources increases only very weakly as the off-

TABLE 5  
LIST OF POINT SOURCES DETECTED BETWEEN 0.5 AND 2 keV AND 2 AND 10 keV  
IN THE 3C 273 FIELD—ACIS-S3 CENTRAL CHIP

Row	R.A. (2000)	Decl. (2000)	Net Counts	Background Counts	Flux ( $10^{-15}$ cgs)	$\theta_{\text{off}}$ (arcmin)
0.5–2 keV						
1 <sup>a</sup> .....	12 29 07.0	+02 03 07.9	1095 ± 109	402 ± 0.22	1336 ± 12.9	1.1
2 .....	12 28 59.5	+02 10 50.9	119.71 ± 11.84	7.64 ± 0.03	14.60 ± 1.44	7.1
3 .....	12 29 15.4	+02 05 29.5	99.75 ± 10.35	4.12 ± 0.02	12.17 ± 1.26	3.4
4 .....	12 29 07.2	+02 04 00.8	31.52 ± 6.32	9.05 ± 0.03	3.85 ± 0.77	0.9
5 .....	12 29 08.5	+02 05 53.9	30.10 ± 5.72	2.25 ± 0.02	3.67 ± 0.70	2.4
6 .....	12 29 14.5	+02 01 21.2	18.06 ± 4.75	3.59 ± 0.02	2.20 ± 0.58	3.8
7 .....	12 28 59.3	+02 05 28.3	17.57 ± 4.60	3.49 ± 0.02	2.14 ± 0.56	1.9
8 .....	12 29 02.2	+02 05 24.9	16.53 ± 4.46	3.42 ± 0.02	2.02 ± 0.54	1.6
9 .....	12 28 52.2	+02 05 13.2	16.49 ± 4.47	2.92 ± 0.02	2.01 ± 0.55	3.1
10 .....	12 29 11.1	+02 05 31.0	15.29 ± 4.30	2.92 ± 0.02	1.87 ± 0.52	2.5
11 .....	12 28 52.2	+02 05 13.2	14.11 ± 4.24	2.70 ± 0.02	1.72 ± 0.52	5.2
12 .....	12 29 17.4	+02 08 31.6	11.81 ± 4.15	3.99 ± 0.02	1.44 ± 0.51	5.9
13 .....	12 29 18.0	+02 06 10.8	9.93 ± 3.62	2.51 ± 0.02	1.21 ± 0.44	4.3
2–10 keV						
1 <sup>a</sup> .....	12 29 07.5	+02 03 08.1	4299 ± 68.1	251 ± 0.17	4084 ± 63.4	...
4 .....	12 29 07.3	+02 04 01.7	60.18 ± 8.88	20.01 ± 0.05	57.17 ± 8.44	...
3 .....	12 29 15.4	+02 05 30.0	35.04 ± 6.35	3.98 ± 0.02	33.28 ± 6.03	...
	12 29 11.7	+02 03 13.3	26.98 ± 6.30	12.24 ± 0.04	25.63 ± 5.98	...
2 .....	12 28 58.6	+02 10 53.1	10.55 ± 3.95	3.37 ± 0.02	10.02 ± 3.75	...

NOTE.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds. Position uncertainties are estimated to be within  $\sim 2''$ .

<sup>a</sup> Source (3C 273) affected by pile-up, so values should be regarded as only indicative.

axis distance increases. This is especially true if one considers only the central chip S3 (where the off-axis distance is confined to  $\lesssim 7'$ ).

At a flux limit of  $1.5 \times 10^{-15}$  ergs cm $^{-2}$  s $^{-1}$ , a source gives  $\sim 15$  and 10 net counts per detected source in RX J0030 and 3C 295, compared with a background of  $\sim 3$  and 1 counts per source extraction area, respectively. Hence, most of the sources should have been detected regardless of position in the FOV. Given the negligible background, most of the sources have a signal-to-noise ratio (S/N) higher than 4 at this flux limit, and we therefore avoid complications caused by “Eddington bias,” as shown by Schmitt & Maccacaro (1986).

The main results of the present study are given in Tables 6 and 7 and are shown in Figures 2 and 3. The table reports the source densities per deg $^2$  measured in the central chip

(S3) of RX J0030 and 3C 295 for two different flux limits ( $1.5$  and  $3 \times 10^{-15}$  ergs cm $^{-2}$  s $^{-1}$ ) for the 0.5–2 keV detections and for  $2 \times 10^{-14}$  ergs cm $^{-2}$  s $^{-1}$  for the 2–10 keV detections. For comparison, we also report in the same table the log  $N$ -log  $S$  obtained from the *ROSAT* Lockman Hole deep field (Hasinger et al. 1998), two recent *Chandra* deep fields (Mushotzky et al. 2000; Giacconi et al. 2000), and from the two comparison fields. The log  $N$ -log  $S$  curves (upper panel of Fig. 2) show the Table 6 numbers plus those obtained at fluxes of  $2.5$  and  $4 \times 10^{-15}$  ergs cm $^{-2}$  s $^{-1}$ . We used the geometric area of one chip (64 arcmin $^2$ ) for RX J0030, 3C 295, and 3C 273, and four chips (256 arcmin $^2$ ) for the anti-Leonid field. These are upper limits on the real area given the uncertainties in instrumental effects (see Appendix B).

It is clear from Table 6 and Figure 2 that the present

TABLE 6  
NUMBER OF SOURCES DETECTED BETWEEN 0.5 AND 2 keV

RX J0030(S3)	3C 295(S3)	<i>ROSAT</i> <sup>a</sup> / <i>Chandra</i> <sup>b</sup> / <i>Chandra</i> <sup>c</sup>	Anti-Leonid (10–3) <sup>d</sup>	3C 273(S3)
$F_{\text{limit}} = 3 \times 10^{-15}$ ergs cm $^{-2}$ s $^{-1}$				
13 ± 3.6	13 ± 3.6	...	7.7 ± 1.4	5 ± 2.2
731 ± 202	731 ± 202	336 ± 31/335/288	436 ± 78	280 ± 126
$F_{\text{limit}} = 1.5 \times 10^{-15}$ ergs cm $^{-2}$ s $^{-1}$				
21 ± 4.6	16 ± 4	...	10.5 ± 1.6	11 ± 3.3
1181 ± 259	900 ± 225	600 ± 70/544/502	578 ± 64	619 ± 187

NOTE.—Errors on the *Chandra* numbers are  $\sqrt{N}$ . For each flux limit, the first row is in chip $^{-1}$  and the second is in deg $^{-2}$ .

<sup>a</sup> *ROSAT* PSPC values and statistical errors obtained from Hasinger et al. 1998.

<sup>b</sup> *Chandra* “best-fit” values obtained from Mushotzky et al. 2000.

<sup>c</sup> *Chandra* “best-fit” values obtained from Giacconi et al. 2000.

<sup>d</sup> Sources were detected in the 4 ACIS-I chips and normalized for 1 single chip (for comparison with RX J0030, 3C 295, and 3C 273).

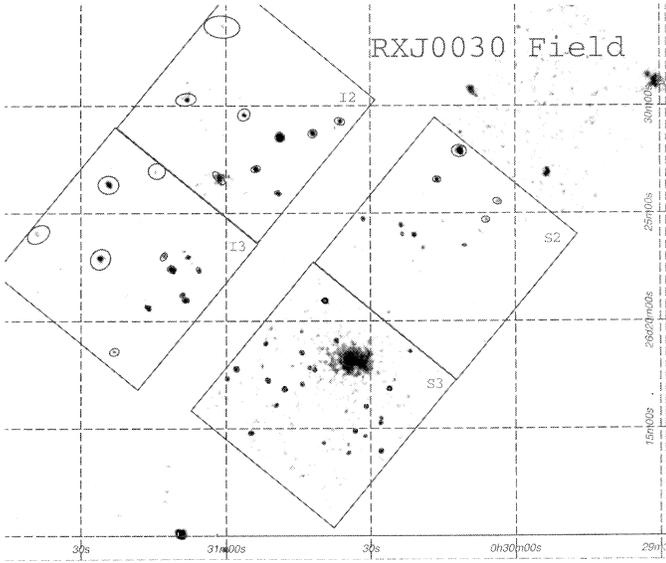


FIG. 1a

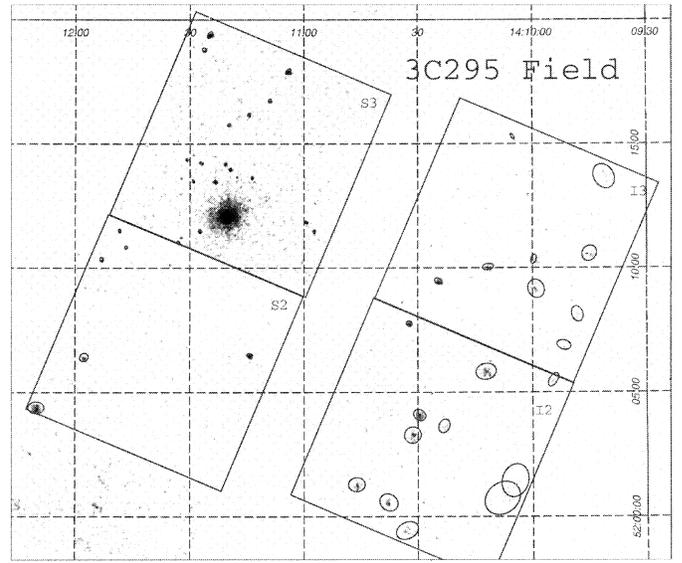


FIG. 1b

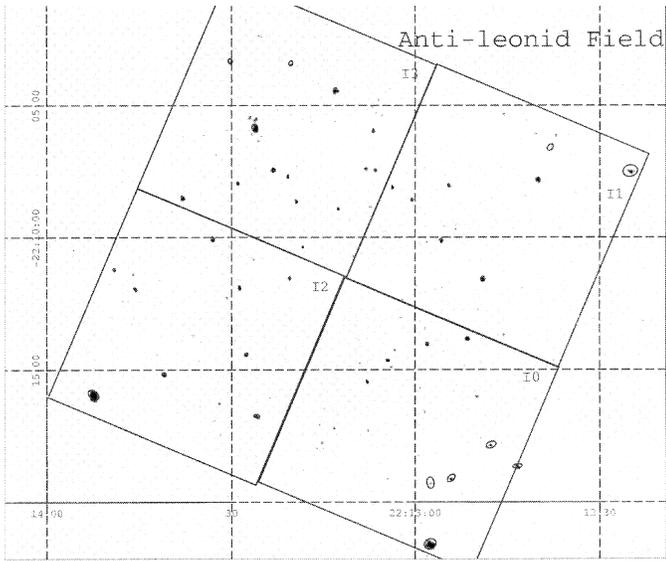


FIG. 1c

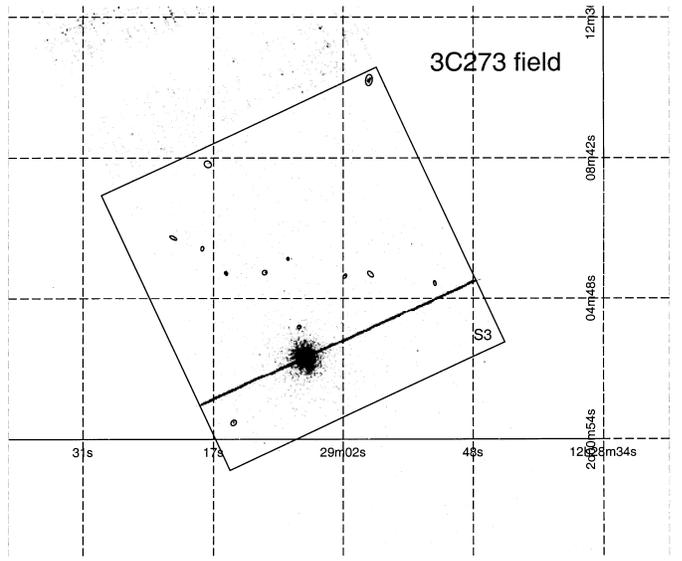


FIG. 1d

FIG. 1.—Images of the (a) RX J0030, (b) 3C 295, (c) anti-Leonid, and (d) 3C 273 fields between 0.5 and 2 keV. The images (north is at the top and east is right) were binned by a factor of 4 (i.e., 1 pixel  $\approx 2''$ ) and gaussian smoothed with a  $\sigma = 1$  pixel. The sources detected as serendipitous sources (those listed in Tables 2, 3, 4, and 5) are marked with elliptical regions. Their sizes are set by deriving the  $1\sigma$  principal axes (and rotation angle) of the source counts distribution for each source and multiplying these by 10 for greater visual clarity. Chip boundaries are marked by the line. Only the chips outlined were used in the analysis (see Appendix A). The line of emission in the 3C 273 field is caused by exposure during read out, which is known to occur with very bright sources.

TABLE 7  
NUMBER OF SOURCES DETECTED BETWEEN 2 AND 10 keV

RX J0030(S3)	3C 295(S3)	ASCA/Model <sup>a</sup>	Anti-Leonid (I0–3) <sup>b</sup>	3C 273(S3)
$F_{\text{limit}} = 2 \times 10^{-14} \text{ ergs cm}^{-2} \text{ s}^{-1}$				
$4 \pm 2^c$	$4 \pm 2$	...	$1.2 \pm 0.5$	$4 \pm 2$
$225 \pm 112^c$	$225 \pm 112$	120/120	$70 \pm 31$	$225 \pm 112$

<sup>a</sup> Linear extrapolation of measured ASCA values from Della Ceca et al. 2000 and expected value from theoretical model of Comastri et al. 1999.

<sup>b</sup> Sources were detected in the 4 ACIS-I chips and normalized for 1 single chip (for comparison with RX J0030, 3C 295, and 3C 273).

<sup>c</sup> First row in  $\text{chip}^{-1}$ ; second row in  $\text{deg}^{-2}$

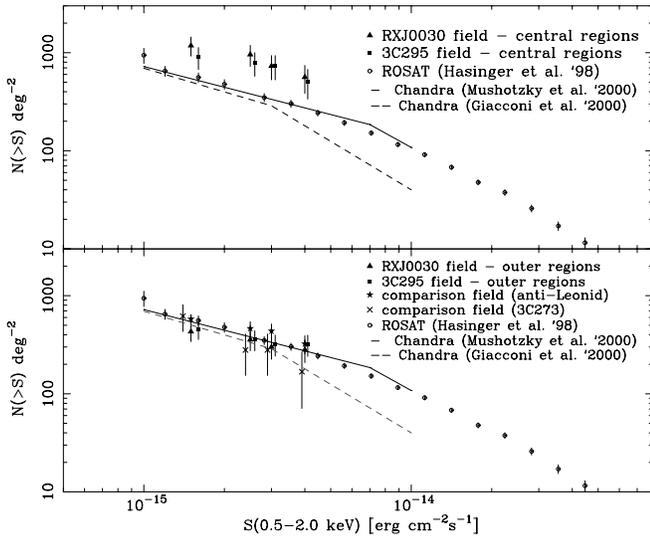


FIG. 2.—*Upper panel:* *Chandra* log  $N$ -log  $S$  between 0.5 and 2 keV measured in the RX J0030 and 3C 295 fields (only central chips, see text for details) compared with the 0.5–2 keV log  $N$ -log  $S$  measured by *ROSAT* (PSPC+HRI) from Hasinger et al. (1998) and *Chandra* “best-fit” results from Mushotzky et al. (2000) and Giacconi et al. (2000). In the Giacconi et al. log  $N$ -log  $S$ , the lack of sources for fluxes between  $3 \times 10^{-15}$  and  $10^{-14}$  ergs  $s^{-1} cm^{-2}$  is attributed to cosmic variance (Giacconi et al. 2000). Reported fluxes for 3C 295 were increased by  $1 \times 10^{-16}$  ergs  $cm^{-2} s^{-1}$  for plotting clarity. *Lower panel:* *Chandra* log  $N$ -log  $S$  between 0.5 and 2 keV measured in the RX J0030 and 3C 295 fields (only outer chips) and log  $N$ -log  $S$  obtained in the two comparison fields (anti-Leonid: whole FOV; 3C 273: chip S3). *ROSAT* and *Chandra* deep-field data points are again plotted for comparison. Reported fluxes for 3C 273 were decreased by  $10^{-16}$  ergs  $cm^{-2} s^{-1}$  for plotting clarity.

0.5–2 keV results strongly suggest an excess number of serendipitous sources in the RX J0030 and 3C 295 fields when compared with *ROSAT* and *Chandra* deep-field counts. Assuming the *ROSAT* value as the true mean source density, the probability of finding a number of sources equal or greater than that observed in RX J0030 is 1% and 0.4% for fluxes of 3 and  $1.5 \times 10^{-15}$  ergs  $cm^{-2} s^{-1}$ , respectively.

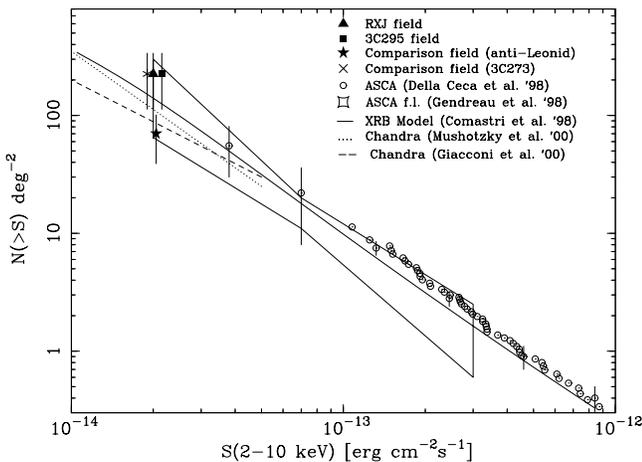


FIG. 3.—*Chandra* log  $N$ -log  $S$  between 2 and 10 keV measured for a flux limit of  $2 \times 10^{-14}$  ergs  $cm^{-2} s^{-1}$  in the RX J0030, 3C 295 (only central chips, see text for details), 3C 273 (S3 chip), and anti-Leonid fields (whole FOV, 4 chips). *ASCA* points (Cagnoni, Della Ceca, & Maccacaro 1998; Della Ceca et al. 2000), *ASCA* fluctuation limits (Gendreau, Barcons, & Fabian 1998), an AGN model for the XRB (Comastri et al. 1995), and *Chandra* deep-field “best-fit” results (calculated between  $1$  and  $5 \times 10^{-14}$  ergs  $cm^{-2} s^{-1}$  from Mushotzky et al. 2000; Giacconi et al. 2000) are plotted for comparison.

For 3C 295, the probabilities are 1% and 8% for the same fluxes.

If the allowed  $1 \sigma$  upper envelope on the *ROSAT* measurements are used instead, the probabilities are about 2 times higher. At a higher flux limit of  $2.5 \times 10^{-15}$  ergs  $cm^{-2} s^{-1}$ , the probability is even lower for RX J0030. Based on measurements of the angular correlation function of X-ray sources derived from a *ROSAT* PSPC survey by Vikhlinin & Forman (1995), we estimate that the rms fluctuations of the X-ray source density are on average  $\sim 20\%$ – $35\%$ . Hence, even if cosmic variance is taken into account (see § 6.1.1), the probabilities remain similar. In summary, the excess is significant at the  $\sim 2 \sigma$  level in both cases, even when the uncertainties in the *ROSAT* estimates are considered. If the two fields are combined, then the probabilities at  $1.5$  and  $3 \times 10^{-15}$  ergs  $cm^{-2} s^{-1}$  become 0.1% and 0.04%, i.e., a significance of  $\sim 3$  and  $3.5 \sigma$ , which is larger at higher than lower fluxes.

In the hard energy band, although we reach a factor 2–3 deeper ( $2 \times 10^{-14}$  ergs  $cm^{-2} s^{-1}$ ) than the *ASCA* and *BeppoSAX* log  $N$ -log  $S$  (Giommi, Perri, & Fiore 2000), the statistics are poorer and the source counts from the four *Chandra* fields are consistent with the Comastri et al. (1999) model for the XRB, the *ASCA* fluctuations, and the *Chandra* deep-field observations (see Fig. 3).

Recently, Brandt et al. (2000) identified some of the *Chandra* 2–8 keV sources in the RX J0030 field (see also § 5.2). They detect nine sources in the central chip. Five out of six of our sources are coincident with their detections and the remaining four are below our conservatively chosen threshold. Brandt et al. do not detect our source 2 in Table 2 because they exclude the chip border. At our hard band flux limit of  $2 \times 10^{-14}$  ergs  $cm^{-2} s^{-1}$ , their derived source density is consistent with the one presented here.

#### 4.2. In the Outer Chips and in the Anti-Leonid and 3C 273 Comparison Fields

Following the procedure explained in § 3, we have calculated the 0.5–2 keV log  $N$ -log  $S$  distribution of the sources detected in the three “external” chips of RX J0030 and 3C 295 (two ACIS-I + 1 ACIS-S chip), in the full anti-Leonid FOV (four ACIS-I chips), and in the S3 chip of the 3C 273 field. These are shown in the lower panel of Figure 2. The results for the anti-Leonid field, the 3C 273 field, and the source density in the outer regions of the cluster fields are fully consistent with the *ROSAT* and *Chandra* log  $N$ -log  $S$  (note that the Mushotzky et al. results were obtained from an observation performed with the S3 chip, like the present cluster fields). These agreements are important because they demonstrate that the excess of sources measured in the central S3 chips of RX J0030 and 3C 295 (§ 4.1) are not caused by an instrumental effect.

The possibility that the “surplus” sources are caused by some statistical fluctuations because of the enhanced background near the clusters halos can be rejected. We find no trend for a higher background (computed locally) around the sources nearer to the clusters (see Tables 2 and 3) and a similar effect would be expected in the 3C 273 field (compare Figs. 1a and 1b with Fig. 1d), which is not observed.

#### 4.3. Spatial Distribution

We computed the density of the sources as a function of off-axis distance for all fields. For the two cluster fields, this clearly corresponds to the (angular) distance from the

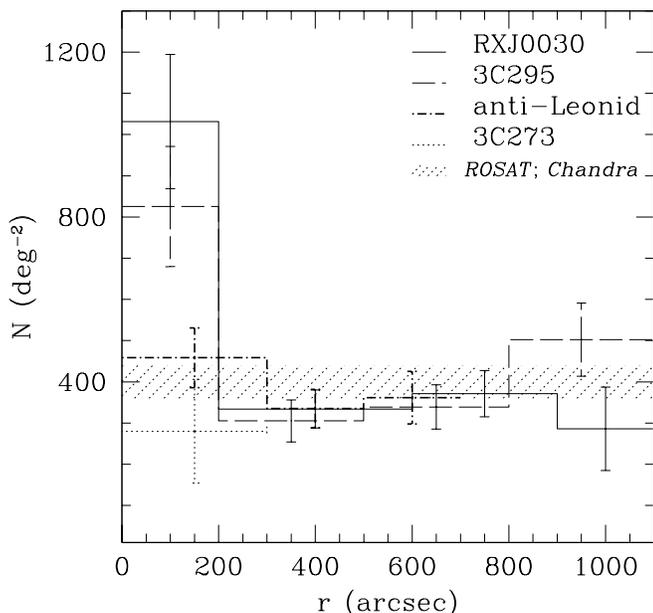


FIG. 4.—Radial distribution of the sources in RX J0030 (solid line), 3C 295 (dashed line), anti-Leonid (thick dash-dotted line), and 3C 273 (dotted line). At  $z = 0.5$ ,  $200''$  correspond to  $\sim 1.4$  Mpc. The dashed area corresponds to the level of the *ROSAT* and *Chandra* log  $N$ -log  $S$  at a flux limit of  $2.5 \times 10^{-15}$  ergs  $\text{cm}^{-2} \text{s}^{-1}$ .

cluster centers since they are almost coincident with the optical axis. In this case, we made a conservative choice for the flux limit of  $\sim 2.5 \times 10^{-15}$  ergs  $\text{cm}^{-2} \text{s}^{-1}$  between 0.5 and 2 keV. We used this value instead of  $1.5$  or  $3 \times 10^{-15}$  ergs  $\text{cm}^{-2} \text{s}^{-1}$  given in Table 6 because we made a trade-off between limiting as much as possible systematic effects at off-axis distance (see § 3) and having the largest number of sources. Starting from the center, bins were constructed by adding  $100''$  to its radius until there were nine sources in each bin. The number of sources per bin were then divided by the area covered by the annuli and the resulting source densities plotted as a function of off-axis distance (Fig. 4). At this flux limit, only five sources were detected in the 3C 273 S3 chip, resulting in a single bin being plotted in Figure 4. This figure suggests, again, that the source densities are higher by a factor  $\sim 2$  near the cluster centers than in the outer regions, while there is no radial dependence in the anti-Leonid comparison field. Moreover, at off-axis distances larger than  $\sim 200''$  (corresponding to  $\sim 1.4$  Mpc at  $z = 0.5$  for  $H_0 = 50 \text{ km s}^{-1} \text{ Mpc}^{-1}$  and  $q_0 = 0.5$ ), the source density in all fields is consistent with the *ROSAT* (Hasinger et al. 1998) and *Chandra* (Mushotzky et al. 2000; Giacconi et al. 2000) log  $N$ -log  $S$ .

## 5. PROPERTIES OF THE SERENDIPITOUS SOURCES

### 5.1. Summed X-Ray Spectra

Here we present the results obtained from the spectral analysis of the co-added spectrum from all the S3 sources in the RX J0030 and 3C 295 central fields (i.e., the average spectra of sources listed in Tables 2 and 3). We also present the summed spectra of all serendipitous sources in the two comparison fields. For both fields, data obtained from the brightest of the sources (the first source in Tables 2, 3, 4, and 5) have been excluded from the summed spectra.

A single spectrum has been constructed that includes the sum of all counts extracted from elliptical regions (chosen to match the spatially varying *Chandra* PSF) centered on the

detected sources, with minor and major axis typically between about  $3''$  and  $7''$ , chosen to ensure inclusion of more than 90% of the PSF encircled energy, at all energies and off-axis distances. Pulse invariant (PI) response matrices released in 1999 October were used. We used only the data between 0.5 and 10 keV, where the matrices are best calibrated. The charge transfer inefficiency problems of the FI CCD (for the anti-Leonid field only) are not corrected for; changes can be expected once more calibrations become available. The spectra were fitted using the *Sherpa* fitting and modeling application included in CIAO. We applied  $\chi^2$  statistics with the Gehrels (1986) approximation of errors in the low-counts regime.

A somewhat critical point of this analysis is that spectral deviations caused by energy-dependent vignetting could artificially steepen our averaged spectrum since, as discussed in Appendix B, the vignetting is larger at higher energies. In the cluster fields (of which only the central chips were considered for the spectral analysis) and in the 3C 273 field, the vignetting should not significantly affect our results since most of the effective area (and therefore counts) is at  $E \lesssim 4$  keV where the vignetting is negligible within  $\sim 5'$ . Indeed the spectra constructed from only the off-axis sources (at greater than  $5'$  off axis) are consistent with that of the on-axis sources. In the case of the anti-Leonid field, we have limited the spectral analysis to the 24 sources detected within an off-axis distance  $\lesssim 7'$  for comparison with the other fields.

Background contributes a significant fraction ( $> 50\%$ ) of the counts at  $E > 4$  keV. We extracted background from several (more than three) large circles of radii  $\lesssim 70''$ , chosen randomly in regions with no detected point sources, and co-added their spectra. We also considered a background chosen from annular regions around each of the source regions, and again added all the counts into a single spectrum. We found that the two background choices gave best-fit parameters consistent with each other to within  $\sim 10\%$ . Hereafter, we report the results obtained with the background determined from the large circles since this has better statistics.

Best-fit results are given in Table 8 and are shown in Figure 5. The main result of this analysis is that the summed spectra are consistent with a single power-law model ( $\Gamma \sim 1.7$ ) with no absorption in excess of the Galactic value. The spectra in the RX J0030 and 3C 295 fields ( $\Gamma \simeq 1.7 \pm 0.2$ ) are slightly flatter than the summed spectrum in the anti-Leonid field ( $\Gamma \simeq 2.3 \pm 0.2$ ) but softer than in the 3C 273 field ( $\Gamma \simeq 1.2 \pm 0.3$ ). The steeper spectrum from the anti-Leonid field could be because of a CTI effect that mostly affects the FI chips.

TABLE 8

BEST FITS OF SUMMED<sup>a</sup> SPECTRA—SINGLE POWER-LAW MODELS

Field	$N_{\text{H}} \equiv N_{\text{Hgal}}$ ( $\times 10^{20} \text{ cm}^{-2}$ )	$\Gamma$	$\chi^2$ (dof)
RX J0030 .....	3.9	$1.72^{+0.13}_{-0.12}$	27.5 (39)
3C 295 .....	1.33	$1.79^{+0.15}_{-0.13}$	21.3 (30)
Anti-Leonid .....	2.45	$2.28^{+0.17}_{-0.15}$	34.7 (42)
3C 273 .....	1.80	$1.24^{+0.25}_{-0.20}$	32.9 (36)

NOTE.—Intervals are at 90% confidence for one interesting parameter.

<sup>a</sup> Computed with 22 sources in RX J0030, 16 in 3C 295, 24 in the anti-Leonid fields, and 12 in 3C 273 (see § 5.1).

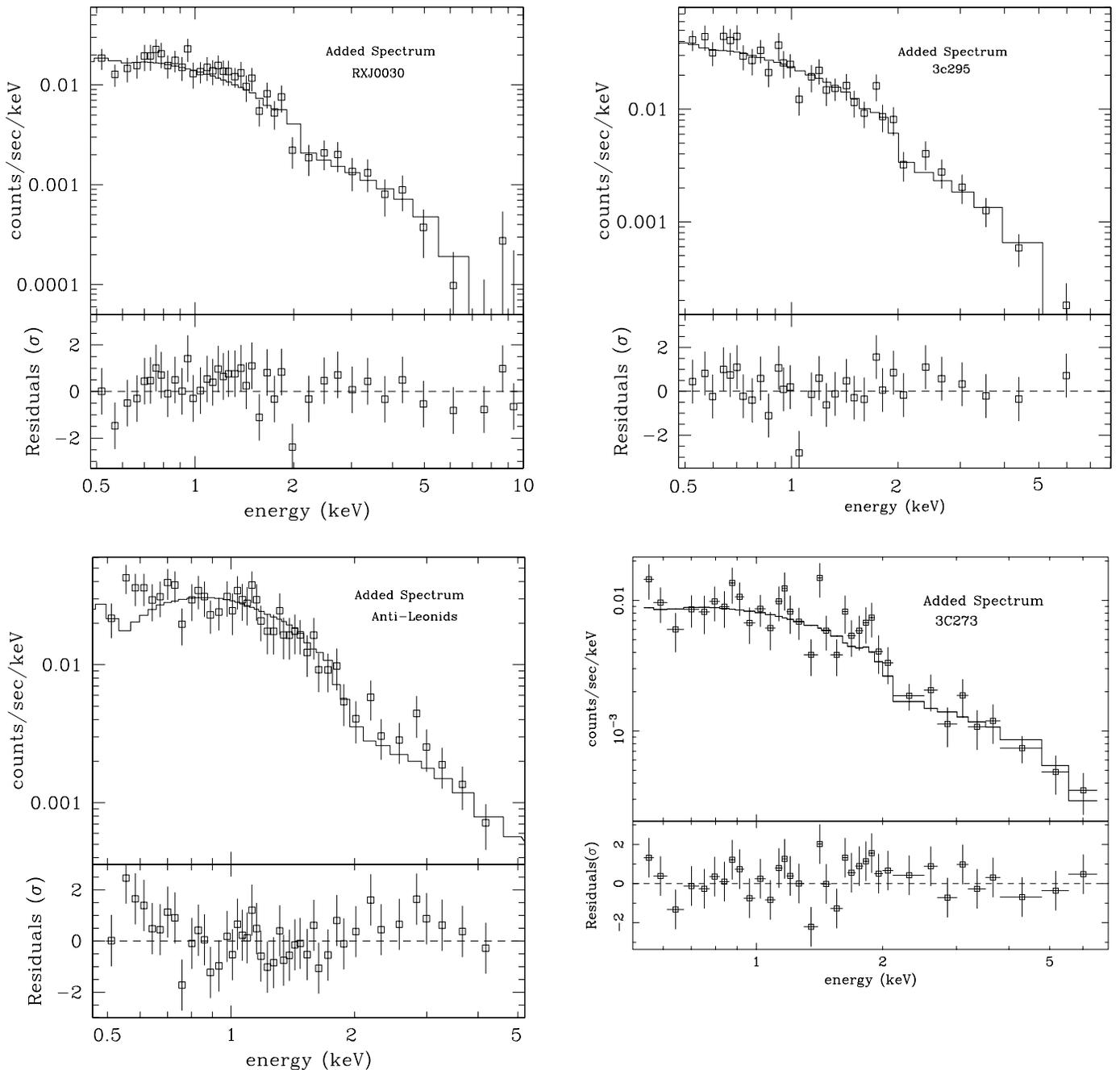


FIG. 5.—Added spectra of all the serendipitous sources (cluster and/or brightest sources excluded) in the RX J0030, 3C 295, anti-Leonid, and 3C 273 fields: best-fit spectra modeled with a single power law as given in Table 8. Data were binned in order to have an S/N > 3 in each energy bin.

The spectra can also be described by a high-temperature ( $kT \sim 2\text{--}4$  keV if  $z = 0$  and  $kT \sim 4\text{--}6$  keV if  $z = 0.5$ ) thermal model with poorly constrained abundances ( $\sim 0.1\text{--}1$  solar) and with  $\chi^2$ -values comparable to the single power-law model. The two models are indistinguishable, which is not surprising given the limited statistics. In addition, we note that the summed spectra of sources detected only in the soft X-ray band are in all cases consistent with the “total” summed spectra, in agreement with the fact that we detect no “hard X-ray only” sources (§ 4.1).

### 5.2. Optical Identifications

We searched for optical counterparts to the S3 X-ray sources in RX J0030 and 3C 295 from the USNO-A2.0

catalog (Monet et al. 1998), which has  $1\sigma$  positional errors  $\sim 0''.25$  (Deutsch 1999) and reaches a  $B$  magnitude limit of about 20. We used a search radius of  $3''$ . For larger radii, the fraction of random matches exceeds 10%. The USNO-A2.0 catalog includes  $B_J$  and  $R$  magnitudes. For ease of comparison with other works, we convert optical magnitudes originally in the  $B_J$  band to the  $B$  band using  $B = B_J + 0.28$  ( $B - V$ ) (Blair & Gilmore 1982), assuming  $(B - V) = (B - R)/2$ . We apply a correction for extinction using the results of Burstein & Heiles (1978, 1982) and assuming  $A_B = 4.0 E(B - V)$ . In addition,  $B$  magnitudes, redshifts, and classifications were obtained from the NED database for two objects in RX J0030 and for five in 3C 295. We also included the recent identifications by Brandt et al. (2000) of

three additional sources in the field of RX J0030. Dressler & Gunn (1992) mapped the central ( $\sim 4' \times 4'$ ) regions of 3C 295 and obtained photometric redshifts for more than 100 galaxies in that area, including two of the *Chandra* source counterparts. In total, we obtain photometric data for 13 sources and redshifts for 10 of these in the RX J0030 and 3C 295 central chips (Tables 2 and 3).

Starting from the  $B$  magnitudes and the derived X-ray (0.5–2 keV) fluxes, we calculate the nominal X-ray to optical slope  $\alpha_{\text{ox}}$  following Stocke et al. (1991).  $B$  magnitudes,  $\alpha_{\text{ox}}$ , redshifts, absolute  $M_B$ , and classification (if available) of the detected sources are reported in Tables 2 and 3.

In RX J0030, four of the five sources with redshifts are foreground galaxies while the bright CRSS QSO has  $z = 0.492$ , which makes it a possible cluster member. However, only five of 23 sources have redshifts. Of the nine sources within  $3'$  of 3C 295, five have redshifts: two are high- $z$  QSOs, two are galaxies at  $z \sim 0.6$ , and one (in the first row in Table 3) is a Seyfert 1 associated with the cluster ( $\Delta z = 0.01$ ). The redshifts of the two galaxies have been estimated from optical colors and have relatively large errors ( $\sim 0.1$ ; Thimm et al. 1994), so these two galaxies could be located at the cluster redshift.

Summarizing, one source (of four with redshifts) in RX J0030 and three sources (of five with redshifts) in 3C 295 are possibly associated with the clusters. Given the current poor statistics and redshift uncertainties, deeper optical imaging and spectroscopy of more of these objects are needed to classify the excess sources and to determine whether they are associated with the clusters.

## 6. DISCUSSION

The *Chandra* observations discussed in this paper strongly suggest a factor  $\sim 2$  overdensity of 0.5–2 keV X-ray sources around two high- $z$  clusters compared to field X-ray sources at a significance level of  $\sim 2 \sigma$  each, or  $3.5 \sigma$  when combined. Given the present statistics, we cannot rule out that we have measured (twice) a statistical fluctuation of the population of field X-ray sources. Clearly, further deeper X-ray observations are required to confirm the reality of the excess. It is notable, however, that this effect has been found around two different clusters (the only two in the *Chandra* archive that are currently public) and at roughly the same significance, making the chance coincidence rather unlikely. Our analysis of the outskirts of the FOVs suggests that the excess of sources disappears at large distance ( $\gtrsim 1.5$  Mpc) from the clusters. At least half of these sources must be the sources that produce the X-ray background (XRB) (possibly a mixture of QSOs and Seyfert-type AGNs—Hasinger et al. 1998; Schmidt et al. 1998; Comastri et al. 1999). If the remaining “surplus” X-ray sources are associated with the clusters, then their average luminosity is  $\sim 10^{42-43}$  ergs  $\text{s}^{-1}$  in either the 0.5–2 keV or 2–10 keV bands. The effect merits some discussion.

### 6.1. On the Nature of the Surplus X-Ray Sources

We shall here consider some possible explanations for the origin of these “surplus” X-ray sources.

#### 6.1.1. Cosmic Variance?

Our view of the universe now includes a weblike network of large-scale structures that include galaxies, clusters, and filaments (Peebles 1993; Peacock 1999). Two-point correlation functions have been able to probe mass clustering on

scales  $\lesssim 5$  Mpc for galaxies (e.g., Small et al. 1999),  $\lesssim 10$  Mpc for QSOs (Shaver 1984; Kundic 1997; La Franca, Andreani, & Cristiani 1998), and  $\lesssim 100$  Mpc for clusters (see Bahcall 1988 for a review). If X-ray sources are distributed like galaxies, then their surface density will have fluctuations caused by the large-scale structure. These fluctuations are known as “cosmic variance.” Could cosmic variance produce the observed effect? Few studies have been made in X-rays (see Barcons et al. 2000 for a recent review). Studies of complete samples of X-ray–selected AGNs that probe the  $\sim 1$  Mpc scale of interest have become possible only in the last few years (Boyle & Mo 1993) and have found positive clustering signals on intermediate scales (0.5–10'—Vikhlinin & Forman 1995; correlation length  $r_c \lesssim 40\text{--}80 h^{-1}$  Mpc, Carrera et al. 1998) and smaller scales ( $r_c \simeq 6 \pm 1.6 h^{-1}$  Mpc—Akylas, Georgantopoulos, & Plionis 2000). The question here is, Is the present overdensity consistent with a random sample of cosmic variance? The answer is no, since the observed amplitude is 20%–30%, i.e., much lower than the factor of  $\sim 2$  fluctuations we measure here. Because we are looking at regions that are centered on distant clusters, we have, however, a highly biased sample. If a filament near a cluster lies mostly normal to the plane of the sky, then a source excess caused by AGNs could be produced near to the clusters. If this is the case, this could well represent the first *direct* measurement of large-scale structure of X-ray–selected sources.

#### 6.1.2. AGNs/Quasars Associated with the Clusters?

Following W. Forman et al. (2000, in preparation), we estimate that the virial radii of RX J0030 and 3C 295, corresponding to a mean gas overdensity of 180, is  $\sim 1$  Mpc (i.e.,  $3'$  at  $z = 0.5$ ). It is possible, therefore, that the sources are physically associated with the clusters.

All the objects with optical counterparts have  $\alpha_{\text{ox}}$  consistent with type 1 AGNs (Elvis et al. 1994). If placed at the clusters' redshifts, all sources, including those with optical counterparts, would have X-ray luminosities of  $L_X(0.5\text{--}2 \text{ keV}) \gtrsim 10^{42}$  ergs  $\text{s}^{-1}$  and  $M_B \lesssim -23$ , again consistent with being bright Seyfert 1 type AGN or low-luminosity QSOs.

Supporting the AGN hypothesis are: (1) the shape of the summed spectra, which are similar to that commonly seen in Seyfert type 1 galaxies, (2) the large luminosities (few  $\times 10^{42}$  ergs  $\text{s}^{-1}$ , Tables 2 and 3), and (3) the  $\alpha_{\text{ox}}$  of the few X-ray sources with optical counterparts (Tables 2 and 3). All three properties are typical of Seyfert galaxies and QSOs, which makes this explanation attractive.

This, however, would be quite strange because AGNs are normally rare in clusters. Although the galaxy environment of AGNs is not well studied (Krolik 1999) in the local universe, AGNs are known to occur more frequently ( $\sim 5\%$ ) in field galaxies than in the nearby ( $z < 0.1$ ) cluster of galaxies ( $\sim 1\%$ , Osterbrock 1960; Dressler, Thompson, & Sackett 1985). There is no strong increase in the number of optically selected AGNs in typical higher  $z$  clusters (e.g., Ellingson et al. 1997, 1998; Dressler et al. 1999).

The clusters 3C 295 and RX J0030 may not be typical high- $z$  clusters. Dressler & Gunn (1983) and Dressler et al. (1999) measured a frequency of AGNs of  $\sim 10\%$  in 3C 295, i.e.,  $\sim 10$  times larger than in other distant clusters (see Table 7 in Dressler et al. 1999). The pioneering X-ray work of Henry et al. (1985) also pointed out that the (X-ray–selected) AGN population of 3C 295 is larger than in low-redshift clusters. In conclusion, the AGN/quasars

hypothesis requires 3C 295 and RX J0030 to be intrinsically unusual.

#### 6.1.3. *Powerful Starburst Galaxies Associated with the Clusters?*

The limited statistics of the spectra of these faint X-ray sources do not allow us to place stringent constraints on the origin of their X-ray emission. The spectra are equally well described by a single power-law model or by a thermal model with  $kT \sim$  few keV (§ 5.1), consistent with the spectra of nearby starburst galaxies (Ptak et al. 1997; Dahlem, Weaver, & Heckman 1998; Cappi et al. 1999). The spatial extents of the X-ray sources, compared with the *Chandra* PSF, are consistent with being pointlike. The most stringent cases (on-axis) place a limit of  $\lesssim 2''$  on the radius of the X-ray-emitting region, which, at  $z = 0.5$ , corresponds to a poorly constraining limit of  $\sim 15$  kpc (starbursts have typical dimensions of  $\lesssim 2$  kpc, e.g., Heckman, Armus, & Miley 1990). However, this hypothesis appears to be rather unlikely because the X-ray luminosities associated with the detected X-ray sources would be a factor of  $\sim 10$ – $100$  times larger than usually observed in nearby starburst and normal galaxies (Fabbiano, Kim, & Trinchieri 1992; David, Jones, & Forman 1992). Moreover, there is no suggestion in the literature for a significant number of starbursts in distant clusters—not any more prevalent than in the field (e.g., Balogh et al. 1999).

Recent *Chandra* deep surveys (Hornschemeier et al. 2000; Mushotzky et al. 2000; Fiore et al. 2000) have nevertheless found X-ray bright galaxies at similar redshifts with luminosities as high as a few  $\times 10^{42}$  ergs  $s^{-1}$  that could either be ellipticals or star-forming galaxies. And, as discussed in the previous section, 3C 295 and RX J0030 could be unusual. Indeed, the 3C 295 cluster was first picked out (Butcher & Oemler 1978) as one of the first clusters discovered to have a larger percentage of blue galaxies than in nearby clusters (“Butcher-Oemler effect,” but see Dressler & Gunn 1983). The colors of the blue galaxies suggest they are probably undergoing star formation; they may indeed be powerful starbursts in spiral galaxies surrounding the cluster (Dressler & Gunn 1983; Poggianti et al. 1999).

#### 6.1.4. *Gravitationally Lensed Sources?*

Most rich clusters produce weakly distorted images of background galaxies/QSOs in the optical band (weak lensing) and occasionally giant arcs (strong lensing). Both 3C 295 (Smail et al. 1997) and RX J0030 (B. McNamara, 2000, private communication) show evidence of such effects, which implies cluster masses of the order of  $\sim 10^{14} M_{\odot}$  each. These large masses are also supported by the large ( $1300 \text{ km s}^{-1}$ ) velocity dispersion in 3C 295 (Dressler & Gunn 1992). Both clusters may, therefore, have large Einstein angles of up to  $\sim 1'$ , depending on the redshifts of the background sources (e.g., Kochanek 1992). One may naively expect an increased number of serendipitous sources near the clusters arising from the amplification of distant (and fainter) background sources. Optical evidence for such an effect has been claimed by Rodrigues-Williams & Hogan (1994), who report a statistically significant overdensity (by a factor 1.7) of high-redshift QSOs in the directions of foreground galaxy clusters. In 3C 295, nine (of 17) X-ray sources are located in an annulus of radii  $\sim 1'$ – $3'$  from the cluster center (see Fig. 1b), and four of these sources (sources numbered 3, 5, 8 and 13, Table 3) have redshifts larger than the cluster’s redshift.

However, the magnification caused by gravitational lensing produces two opposing effects on the source number counts (e.g., Croom & Shanks 1999): individual sources appear brighter by  $\mu$  (where  $\mu$  is the amplification factor), raising the number detected; but the area of sky lensed is reduced by the same factor  $\mu$ . The net effect depends on the slope of the  $\log N$ - $\log S$  curve. At high fluxes where the  $\log N$ - $\log S$  is steep, sources are added to counts; but at faint fluxes, the flatter  $\log N$ - $\log S$  produces a *deficit* in the source counts (Wu 1994). Refregier & Loeb (1997) predict on average a *reduction* in the surface density of faint resolved sources at fluxes fainter than about  $10^{-15}$  ergs  $\text{cm}^{-2} \text{ s}^{-1}$  (0.5–2 keV). An upturn in the  $\log N$ - $\log S$  at very faint fluxes (e.g.,  $\lesssim 2 \times 10^{-16}$  ergs  $\text{cm}^{-2} \text{ s}^{-1}$ ) would be needed to create the observed excess.

#### 6.2. “Contamination” of the $L_X$ - $T$ Relation of High- $z$ Clusters

Point sources contribute a substantial fraction of the clusters’ fluxes, at least between 0.5 and 2 keV (see § 3.2). Therefore, if this finding applies to a significant fraction of medium- to high-redshift clusters, past X-ray measurements of luminosities and temperatures of distant clusters obtained with low-resolution experiments (see, e.g., the *ASCA* measurements by Mushotzky & Scharf 1997) have probably been “contaminated” by the surrounding unresolved surplus sources (note that the contribution from field sources is excluded through normal background subtraction). The present data indicate a possible contamination of the cluster’s fluxes of up to  $\sim 40\%$  and  $\sim 15\%$  in RX J0030 and 3C 295, respectively. The effect on the observed  $L_X$ - $T$  relation (e.g., Wu, Xue, & Fang, 1999) will depend on the way in which the contamination varies with redshift, with cluster size, and on the surplus source spectra. The sources have a summed spectrum of  $kT \sim 5$  keV (cluster rest frame); therefore, at the 0<sup>th</sup>-order, these would lower the observed temperature of clusters with  $kT > 5$  keV (i.e., the highest luminosity clusters) but increase that of lower temperature clusters. As a result, these effects should increase the scattering and possibly modify the shape of the true  $L_X$ - $T$  relation of high- $z$  clusters. Any changes to the  $L_X$ - $T$  relation of clusters would have important consequences for using clusters as cosmological probes (e.g., Henry 2000). More studies using *Chandra* and *XMM* are required in order to quantify the implications of such an effect.

## 7. CONCLUSIONS

We have presented a *Chandra* study of an overdensity of X-ray sources around two  $z \sim 0.5$  clusters (RX J0030 and 3C 295). The observed source density near these clusters appears to be a factor of 2 times larger than expected on the basis of the *ROSAT* and *Chandra*  $\log N$ - $\log S$  or by comparison to noncluster *Chandra* fields and to the outskirts of the cluster FOVs. The radial distribution of the sources suggests that they are indeed concentrated within  $\sim 200''$  of the clusters. The effect is significant (at  $\sim 2 \sigma$  per cluster or  $3.5 \sigma$  when combined) in the 0.5–2 keV energy band, but not in the 2–10 keV band where the statistics are too poor. Deeper X-ray observations are needed to confirm this result unambiguously. For both fields, the summed X-ray spectra of the faint sources are consistent with a power-law spectrum with  $\Gamma \sim 1.7$  and no intrinsic absorption. If the sources are at the redshifts of the clusters, their average

luminosity is  $\sim 10^{42-43}$  ergs  $s^{-1}$  making Seyfert-like AGNs/quasars the most likely counterparts. The number of redshifts currently available, however, is too small to be conclusive.

Possible explanations of the apparent overdensity may be a statistical variance of cosmic background sources; AGNs/quasars and/or powerful starburst galaxies associated with the clusters; and gravitationally lensed sources. None of these explanations is without problems, however. Only follow-up X-ray and optical observations (redshift measurements) will determine the true cause.

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## APPENDIX A

### DATA CLEANING AND REDUCTION

The data from the entire FOVs were cleaned and analyzed using the *Chandra* Interactive Analysis of Observations (CIAO) software (release V1.1, M. Elvis et al., 2000, in preparation; see also <http://asc.harvard.edu/cda/>). The data were first filtered to include only the standard event grades 0, 2, 3, 4, and 6 and energies between 0.1 and 10 keV. All hot pixels and bad columns were removed. Time intervals with large background rate (i.e., larger than  $\sim 3\sigma$  over the quiescent value) were removed chip-by-chip, yielding different exposures for each chip. The differences were small,  $\lesssim 5\%$ , for the 3C 295, anti-Leonid, and 3C 273 fields but were  $\sim 25\%$  for RX J0030, where the two BI chips in the RX J0030 field had  $\sim 30$  ks exposure compared with  $\sim 40$  ks for the FI chips. This larger difference is because of the larger and more variable background flux observed in the BI chips, which, being thinner, are less effective in rejecting background high-energy particles (Markevitch 1999).<sup>10</sup> Data from the outer FI S4 chip and the BI S1 chip were excluded from the analysis of RX J0030 and 3C 295 because the effective exposure of S4 was  $\sim 5$  times smaller than for the other chips (indicative of a very noisy background) and because S1 was more than  $10'$  from the optical axis, which reduces its sensitivity by 20%–50% because of vignetting. For the anti-Leonid field, we also excluded the data from the S2 chip because of its large ( $> 10'$ ) off-axis distance. For the 3C 273 field, we considered only the BI S3 chip, which is sufficient for the scope of the present study. We excluded in this field the area underneath the line emission caused by exposure during read-out.

The 3C 273 field provides a comparison for the two cluster fields that is independent of any systematic variations between CCDs, since these three fields were all observed with the same chip, S3. The larger area covered by the four FI chips in the anti-Leonid observation lowers the statistical uncertainty in the noncluster number counts and allows a comparison of the sources' spatial distribution (see § 4.3). The FI chips have significantly lower background flux and a somewhat lower effective area at  $E \lesssim 1$  keV than BI chips. However, these effects should not significantly affect the present results because appropriate matrices (constructed with CIAO) were used to account for the different instrument responses and the background level is, in fact, negligible in both FI and BI at the flux limit used here (§ 4.1).

## APPENDIX B

### THE SOURCE DETECTIONS

In its simplest form, *wavdetect* consists of correlating the image data with wavelet functions in successively larger scale versions of the wavelet, comparing the resulting “correlation maps” to a local background and detecting sources above a given threshold (see footnote 9 for Dobrzycki et al. 1999; Freeman et al. 2000). We considered only the 0.5–2 keV and 2–10 keV energy bands where *Chandra* is best calibrated, and which allows direct comparison with *ROSAT*, *ASCA*, and previous *Chandra* results (see § 4.1). The original data were binned by four, yielding pixels of  $\sim 2''$  on a side, in order to obtain  $\sim 1000 \times 1000$  pixel images (for the whole FOV). The *wavdetect* software was run on several scales in order to match the (variable) dimensions of the PSF over the FOV. Aspects of the detection method include

1. The computation of the correlation maps using an FFT;
2. The computation of a local, exposure-corrected and normalized (i.e., flat-fielded) background in each pixel;
3. Its applicability to the low-counts regime of *Chandra*, as it does not require a minimum number of background counts per pixel for accurate computation of source detection thresholds;
4. Its applicability to multiscale data (i.e., extended sources);

<sup>10</sup> M. Markevitch, 1999, calibration report of November 18, <http://asc.harvard.edu/cal/>.

5. Its applicability to the *Chandra* FOVs (i.e., the algorithm recognizes where the aim point is located);
6. A full error analysis.

As a result, it is substantially more efficient than a sliding-cell technique in detecting weak point sources in crowded fields and extended sources.

We performed several runs of *wavdetect*, setting the probability of erroneously associating a background fluctuation in a pixel with a detection to  $3.4 \times 10^{-6}$ ,  $1.0 \times 10^{-6}$ ,  $2.9 \times 10^{-7}$ , and  $1.0 \times 10^{-7}$ . These probabilities correspond approximately to Gaussian equivalent significances of 4.5, 4.7, 5, and 5.2  $\sigma$ , respectively. ACIS-I simulations and other calculations (Freeman et al. 2000) show that these thresholds produce  $\leq 3$ , 1, 0.3, and 0.1 false detections per four chips FOV, respectively. We set the threshold to correspond to one expected spurious source (probability of  $10^{-6}$ , 4.7  $\sigma$ ).

The *Chandra* dither pattern creates rapid changes in exposure near the edges of the chips that are not currently accounted for in either source-detection algorithm. As a result, this introduces a significant number of spurious detections of extremely faint sources with  $\leq 6$  counts. We therefore introduced an additional selection criterion: that the sources should also have a significance level  $\text{SN} > 3$ , where SN (an output of *wavdetect*) is defined as the number of source counts divided by the Gehrels (1986) standard deviation of the number of background counts, in order to be considered as real. We checked and found that none of the excluded sources had a flux larger than the flux limit used in § 4.1.

We also checked the results from *wavdetect* by comparing them to the *celldetect* (see Dobrzycki et al. 1999 in footnote 9) output. The algorithm *celldetect* is a robust sliding-cell algorithm in CIAO that utilizes a sliding box of variable size according to the position in the FOV and according to a predefined encircled energy (for a PSF calculated at a given energy). The background is calculated, locally, in annuli centered on the cells with an area equal to the source area (typical sizes of the boxes were 6"–10" per side, on-axis, and 10"–15" per side, at  $\leq 5'$  off-axis). We found that, for sources brighter than  $\sim 15$  net counts, the match was better than  $\sim 90\%$  in both the obtained positions and number of source counts, confirming that the *wavelet* algorithm is working well.

## APPENDIX C

### THE SOURCE COUNTS AND FLUXES

Source count rates were obtained using *wavdetect* from regions with typical radii of  $\sim 3''$  (on-axis) and  $10''$  (off-axis). Comparing the *wavdetect* cell regions with the expected dimensions of the PSF at the source positions, we estimate that our extraction regions contain more than 99% of the PSF encircled energy. Background counts were estimated locally for the same area on the basis of the background map produced by *wavdetect* (Freeman et al. 2000).

The measured counts were first corrected for vignetting and then converted to an emitted, unabsorbed flux. The vignetting is larger for photons with  $E > 4$  keV and, at those energies, the effective area is reduced by  $\sim 40\%$  at  $10'$ . We applied a correction for the effective area off-axis using the preflight calibration. However, this correction was typically not large, since our detections are dominated by photons with  $E < 4$  keV (see Fig. 5). The vignetting  $10'$  off-axis drops to  $\sim 15\%$  for photon energies of  $\sim 1.5$  keV, where *Chandra*'s effective area peaks. The correction we applied is based on Figure 4.3 of the "*Chandra* Proposer's Observatory Guide" (1999), which we approximate using a linear function of the form:

$$V_{\text{cor}} = 0.97 + 0.0175 \times \theta_{\text{off}},$$

where  $V_{\text{cor}}$  is the vignetting correction factor (to be multiplied to all the counts) and  $\theta_{\text{off}}$  is the off-axis distance in arcminutes.

For the conversion from counts to flux, a power-law spectrum with photon index  $\Gamma = 2$  and  $\Gamma = 1.7$  was assumed for the sources detected between 0.5 and 2 keV and between 2 and 10 keV, respectively. We assumed Galactic absorption columns of 3.91, 1.33, 2.45, and  $1.8 \times 10^{20} \text{ cm}^{-2}$  along the line of sight of the RX J0030, 3C 295, anti-Leonid, and 3C 273 fields, respectively (Stark et al. 1992). These spectral models were chosen to allow a direct comparison with the *ROSAT* PSPC and *ASCA* results (Hasinger et al. 1998; Della Ceca et al. 2000). These assumed spectra are consistent with the summed spectra obtained in § 4.2. Moreover, in these energy bands, the fluxes are only weakly dependent on the spectral slope and the Galactic  $N_{\text{H}}$  value adopted (e.g., a  $\Delta\Gamma = 0.2$  gives a  $\Delta\text{flux}$  of less than 5%). Systematic instrumental errors in the absolute effective area and energy scale are conservatively below 10%–20% at all the energies considered here (N. Schultz, 2000, private communication). Small changes can be expected once refined calibrations become available.

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