

Particle acceleration in relativistic magnetized collisionless shocks

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SNRs and PWNe in the Chandra Era

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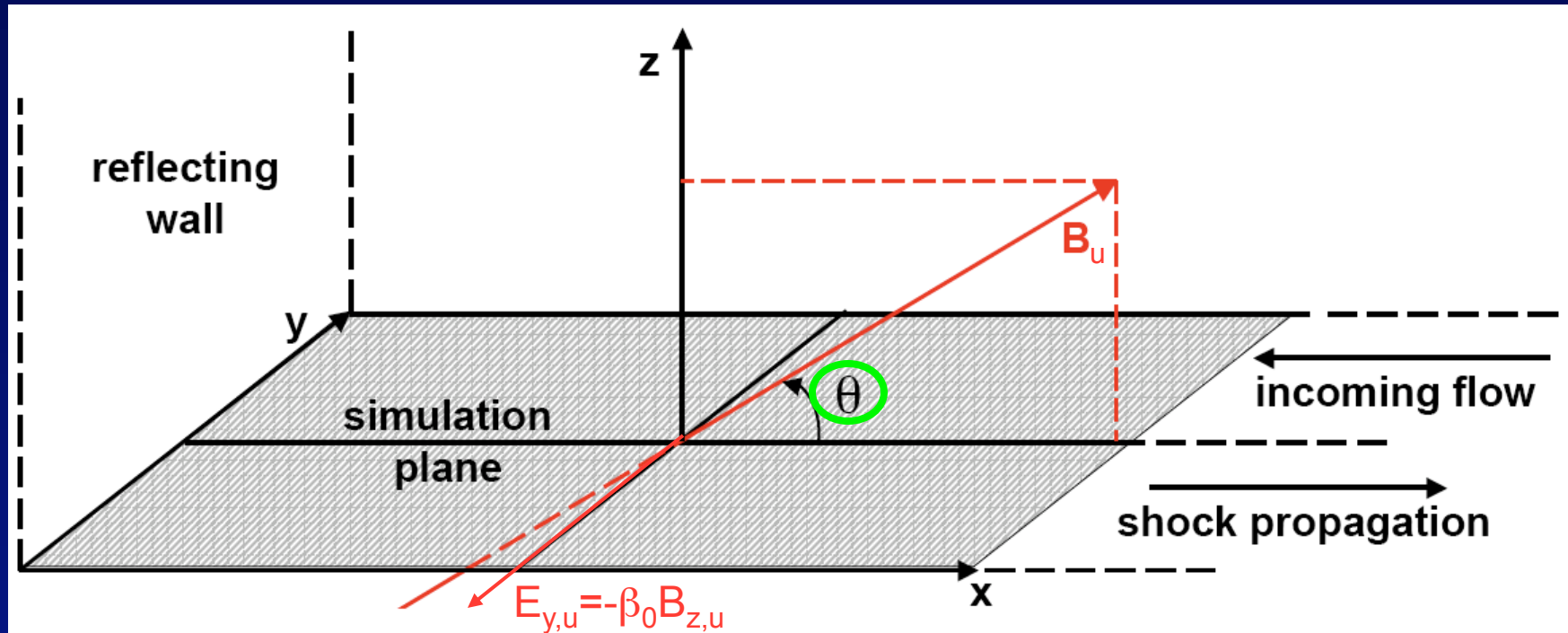
Astrophysical shocks

- In SNRs, GRBs, AGN jets, PWNe: how do collisionless shocks work? How do they produce the observed non-thermal radiation?
- Basic ingredients:
 1. **Magnetic fields**
 2. **Accelerated particles**

How does the efficiency of particle acceleration depend on the magnetic field strength and inclination and the flow composition?

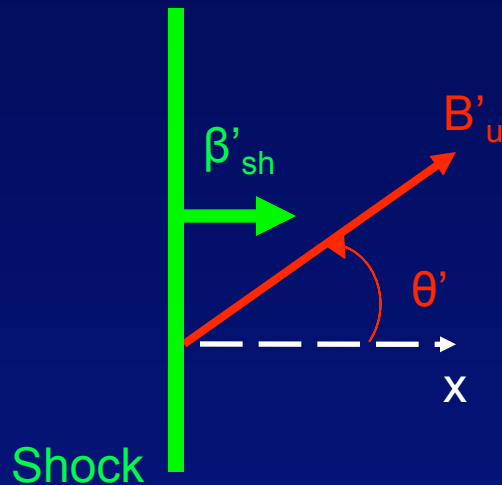
Method: Self-consistent first-principle particle-in-cell (**PIC**) numerical simulations...

Simulation setup



- Injected flow is reflected by a wall; simulation in the **wall** frame
- Upstream flow is e^-e^+ or e^-p^+ ($m_p/m_e=16$) cold plasma with bulk Lorentz factor $\gamma_0=15$ and magnetization $\sigma = B_u^2/(4\pi\gamma_0 n_u m_p c^2) = 0.01-0.1$. We vary the wall-frame magnetic obliquity θ .
- 2.5D simulations ($100 c/\omega_{pe} \times 9000 c/\omega_{pe}$) with out-of-plane magnetic field; main results confirmed by 3D simulations

Subluminal vs superluminal shocks



In the upstream frame

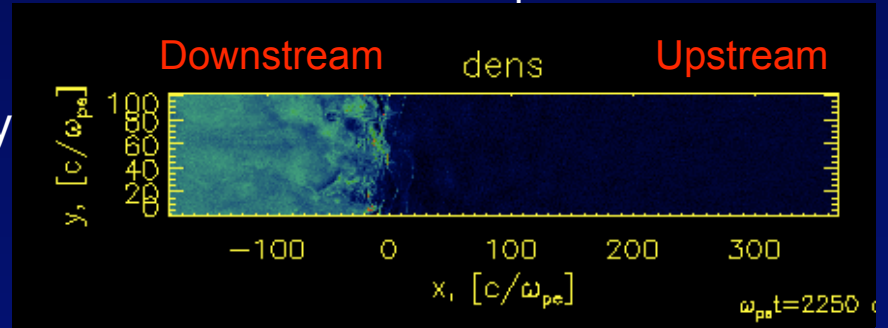
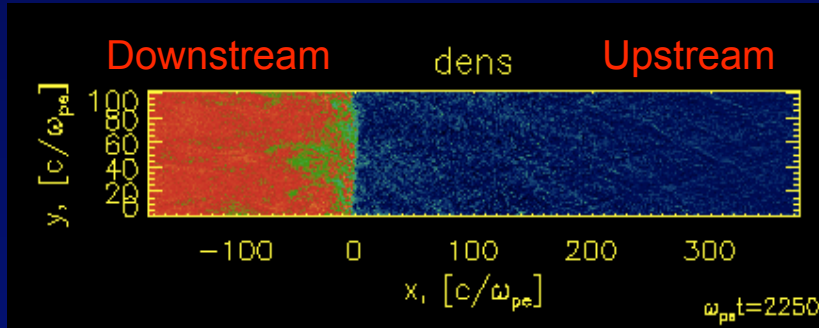
$$\theta'_{\text{crit}} = \arccos(\beta'_{\text{sh}})$$

- In **superluminal** (vs **subluminal**) shocks, a particle sliding along magnetic field lines **CANNOT** (vs **CAN**) return upstream
- For $\gamma_0=15$ and $\sigma=0.1$, the critical obliquity is $\theta_{\text{crit}} \approx 34^\circ$ in the wall frame; in the upstream frame $\theta'_{\text{crit}} \approx 34^\circ/\gamma_0$
- θ_{crit} weakly depends on both γ_0 (≥ 5) and σ ($0.01 < \sigma < 0.3$)
- We expect particle acceleration to be suppressed in superluminal shocks, unless there is strong magnetic turbulence. **But is strong turbulence self-consistently produced by the shock?**

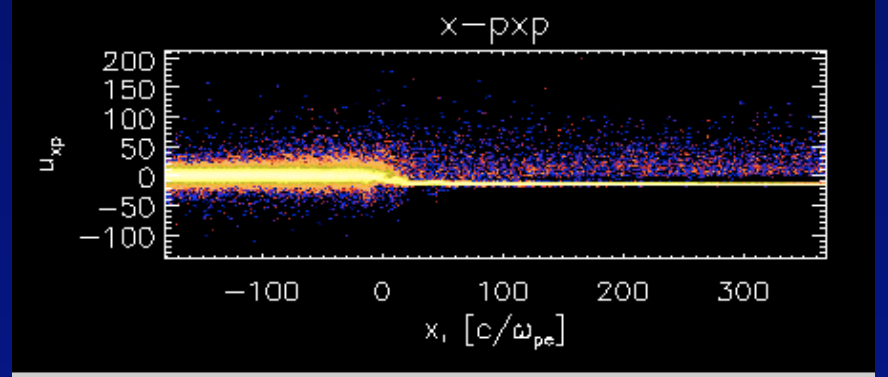
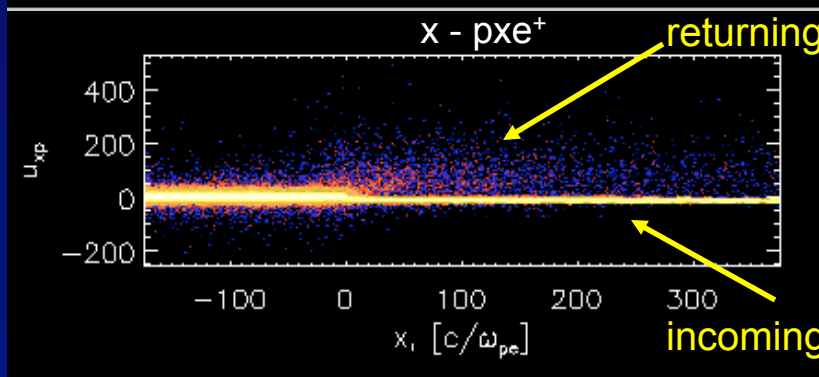
$\theta=15^\circ$: a subluminal shock

e^-e^+ shock

e^-p^+ shock ($m_p/m_e=16$)

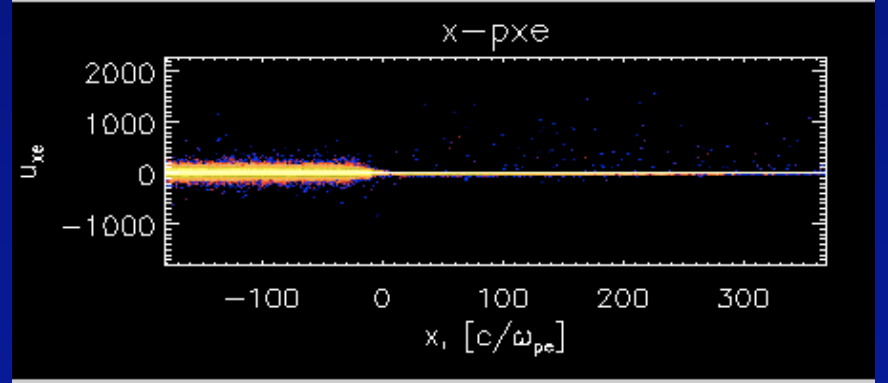
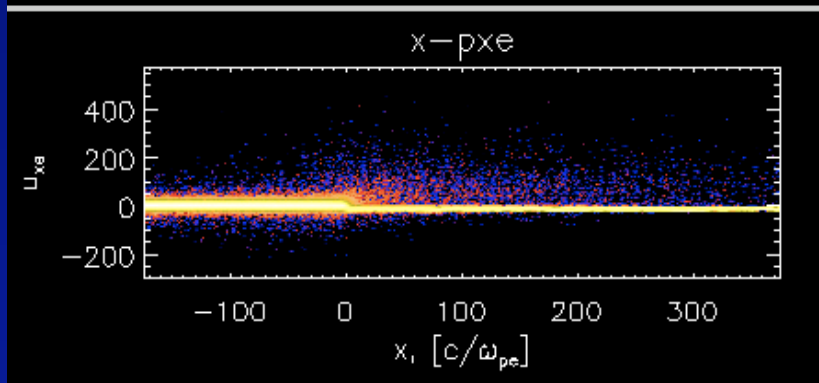


Density



$\gamma_p \beta_{xp}$

$\gamma_e \beta_{xe}$

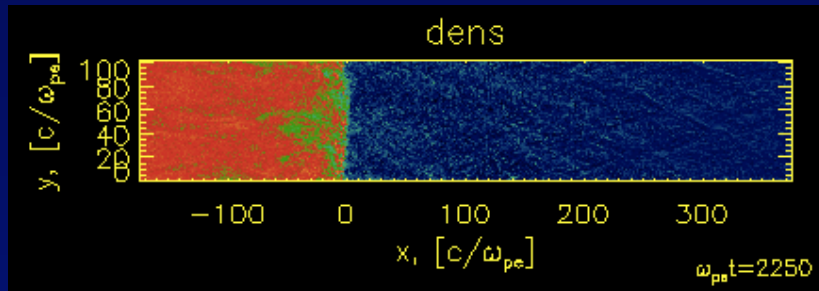


Subluminal \rightarrow Returning particles (mostly **IONS** for e^-p^+ shock)

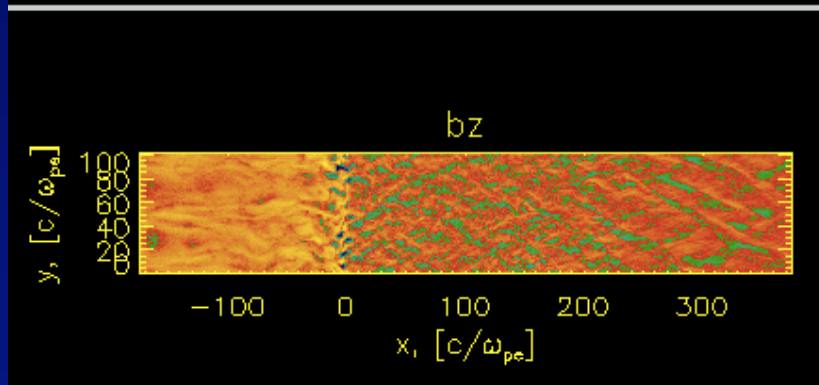
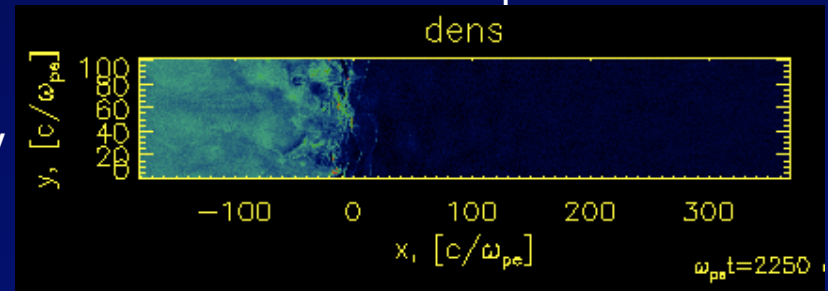
$\theta=15^\circ$: a subluminal shock

e^-e^+ shock

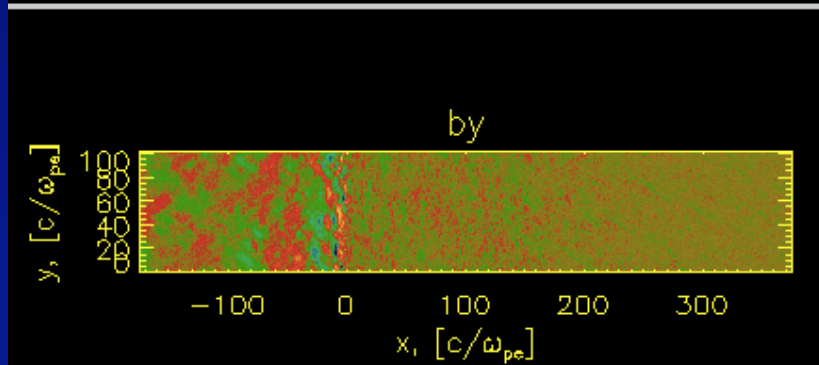
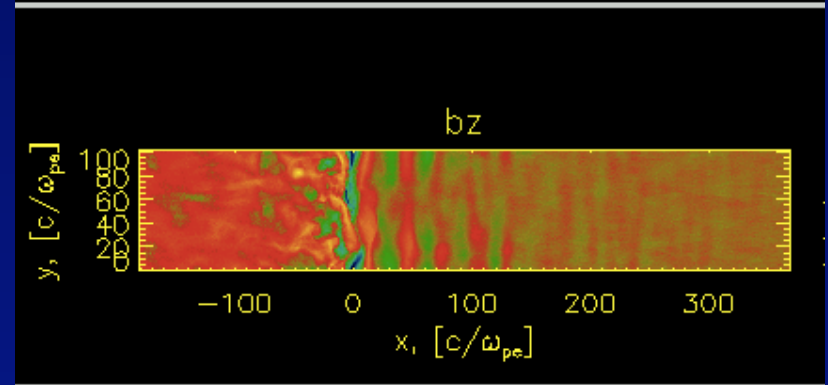
e^-p^+ shock ($m_p/m_e=16$)



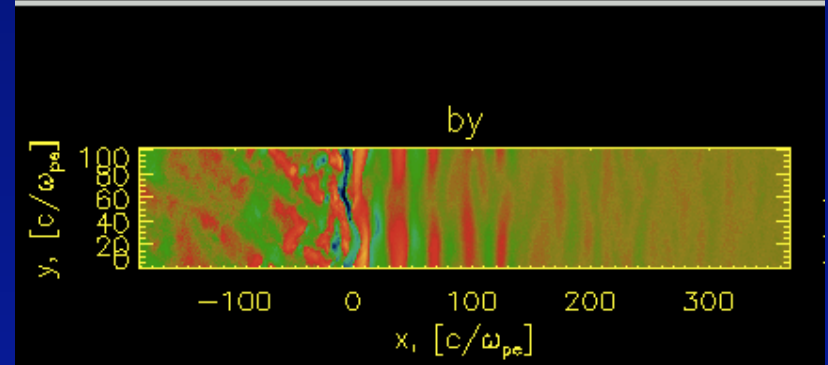
Density



B_z



B_y



Returning particles \rightarrow Upstream waves (oblique vs longitudinal wavevector, linear vs circular polarization)

$\theta=45^\circ$: a superluminal shock

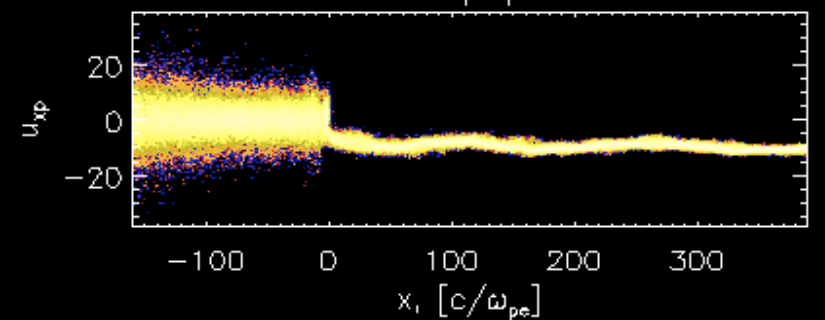
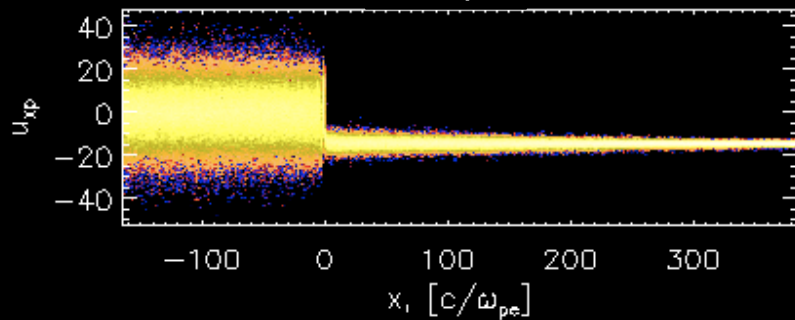
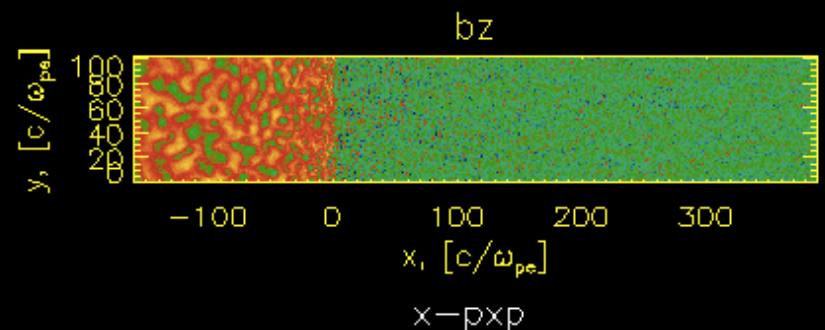
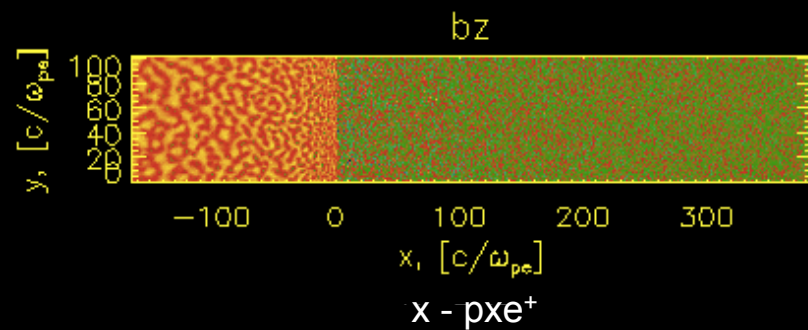
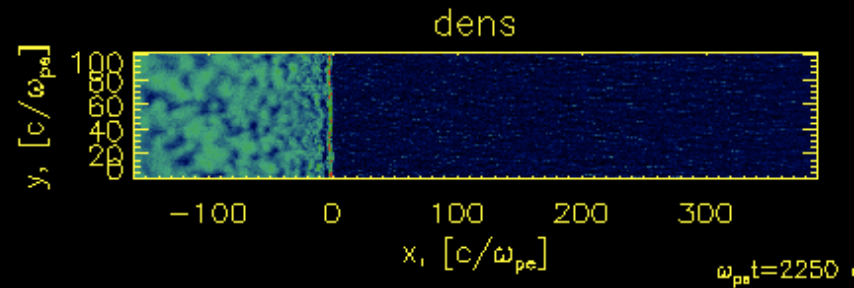
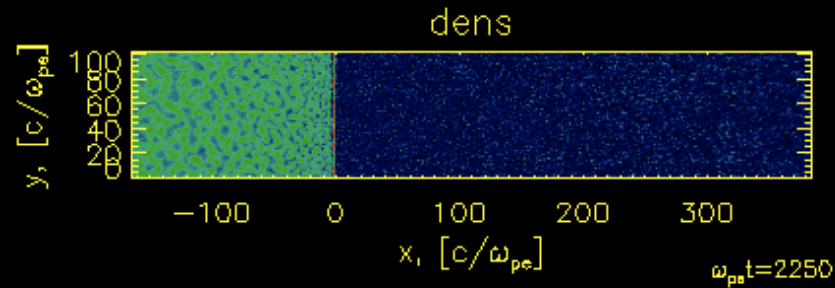
e^-e^+ shock

e^-p^+ shock ($m_p/m_e=16$)

Density

B_z

$\gamma_p \beta_{xp}$



Superluminal \rightarrow No returning particles \rightarrow No upstream waves

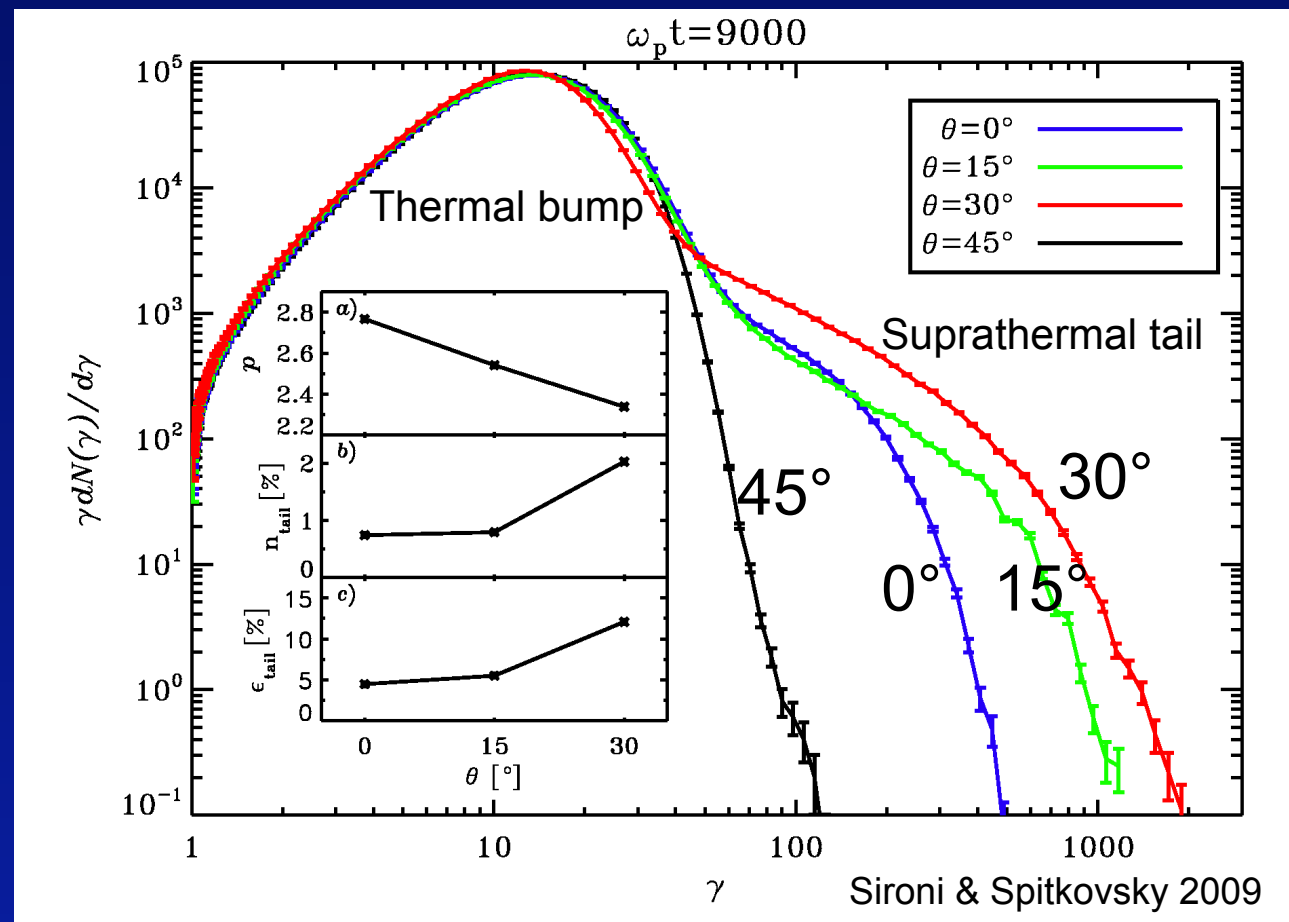
Downstream particle spectra: e⁻-e⁺ shock

- Superluminal shocks DO NOT significantly accelerate, subluminal shocks DO accelerate, the more efficient the closer to $\theta_{\text{crit}} \approx 34^\circ$
- In subluminal shocks, spectra well fitted by 3D low-energy Maxwellian + high-energy power-law tail with exponential cutoff

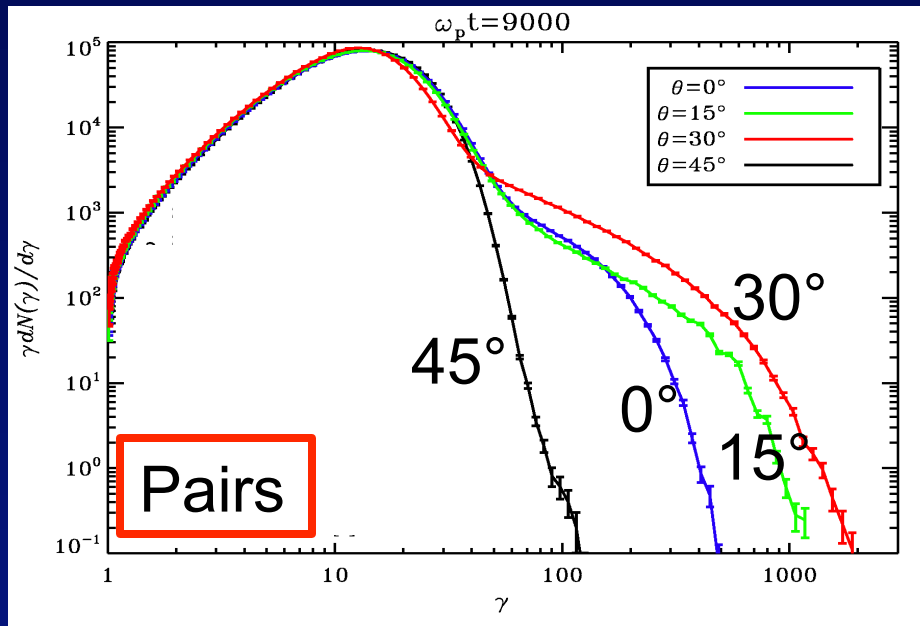
e⁻-e⁺ shock

As θ increases from 0° to 30° :

- Power-law slope from -2.8 to -2.3
- Number fraction from 1% to 2%
- Energy fraction from 5% to 13%

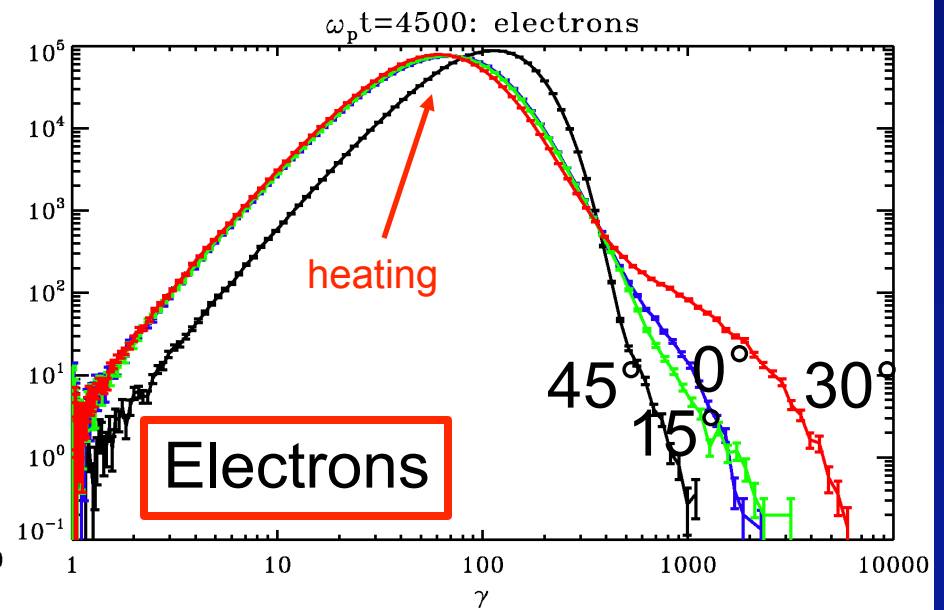
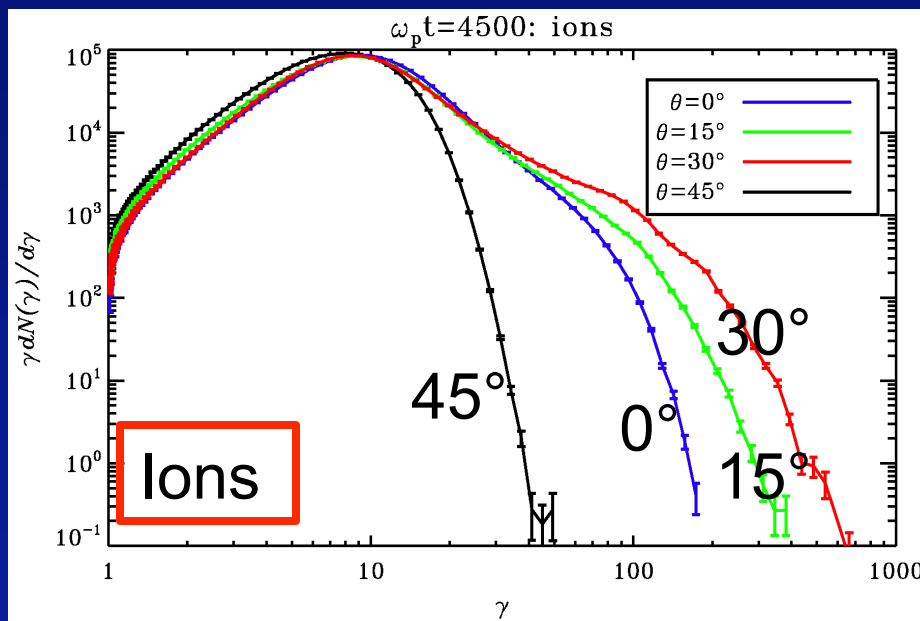


Downstream particle spectra: e-p⁺ shock



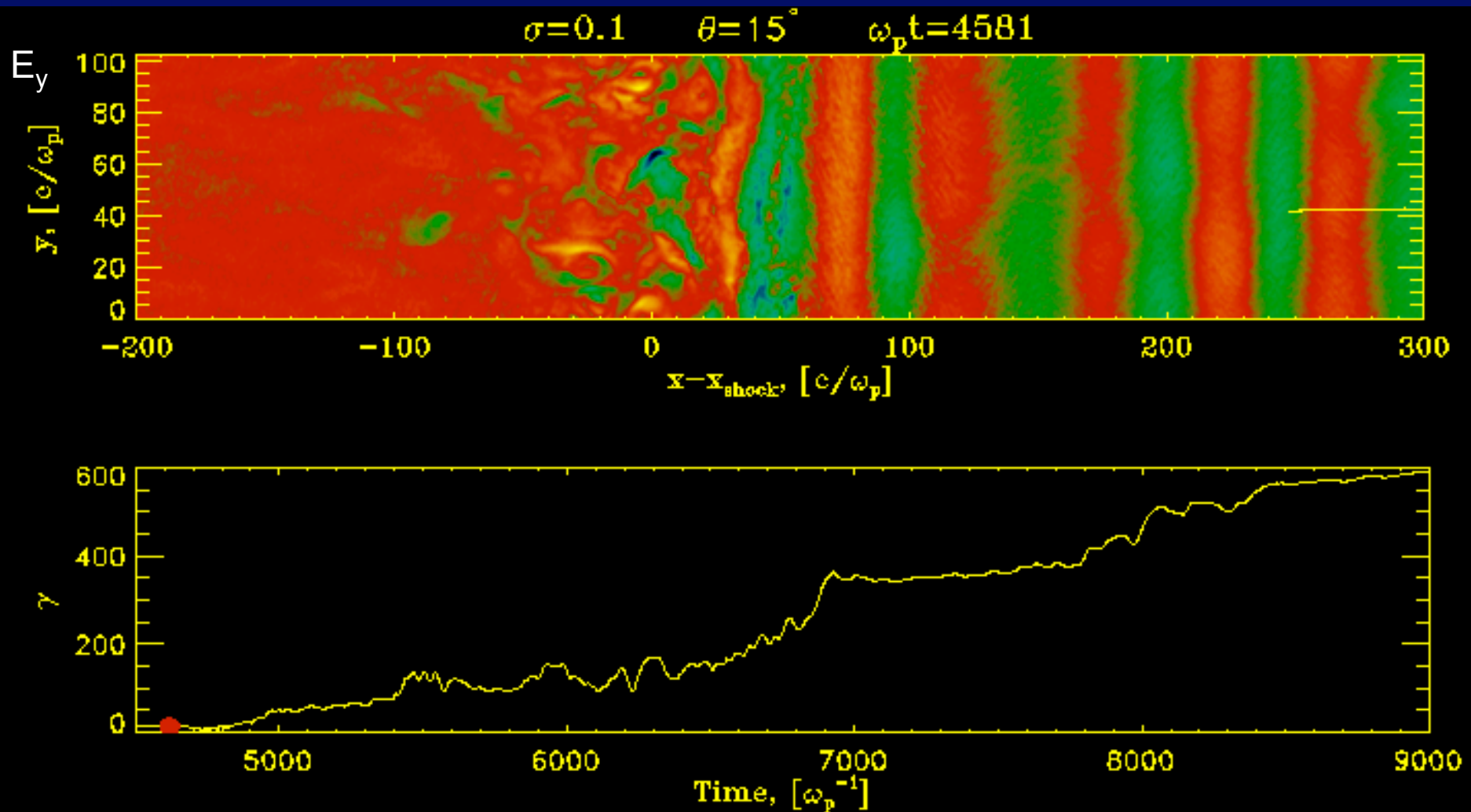
- Ions are accelerated, and their spectra resemble pair spectra in e-e⁺ shocks:
- negligible acceleration in superluminal shocks
 - with increasing θ from 0° to 30°, slope from -3.0 to -2.2, number fraction from 2% to 5%, energy fraction from 10% to 25%

Electron acceleration in e-p⁺ shocks is a factor of 5-10 less efficient than ions



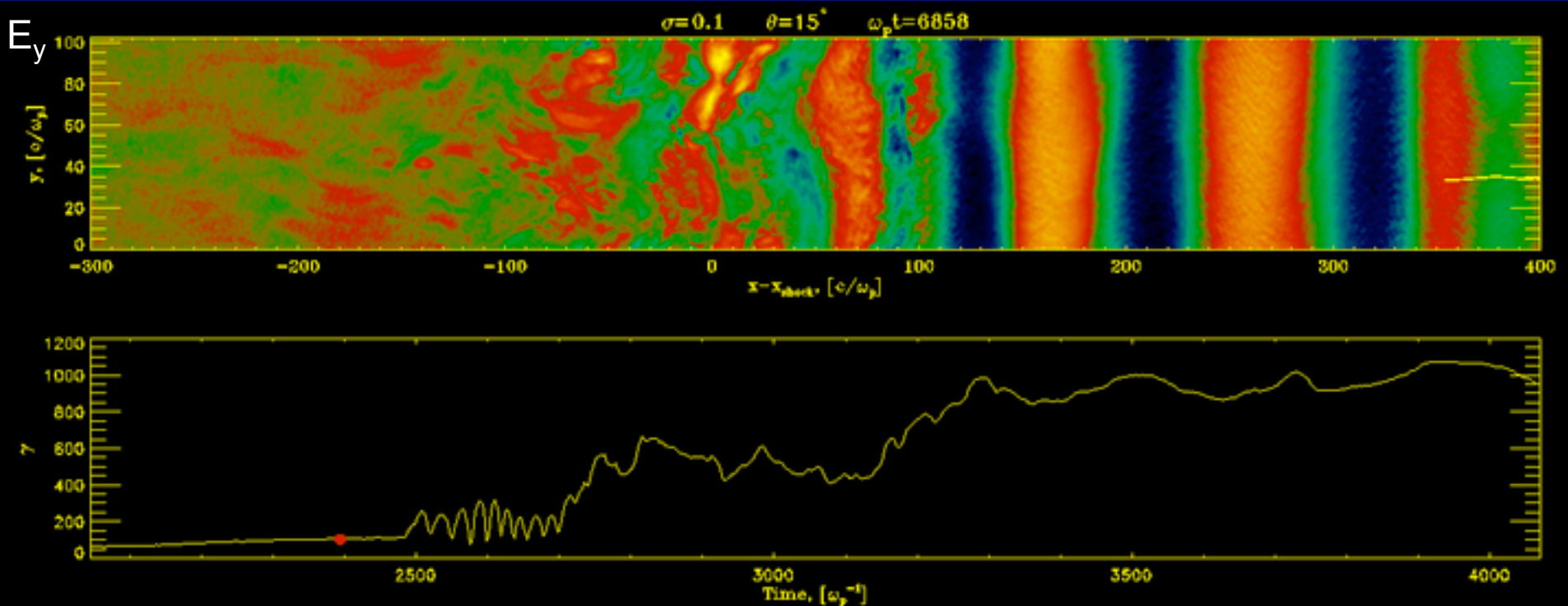
Ion vs electron acceleration (1/2)

$\sigma=0.1$ $\theta=15^\circ$ e⁻-p⁺ shock: **IONS** get accelerated by scattering off the self-generated upstream longitudinal waves

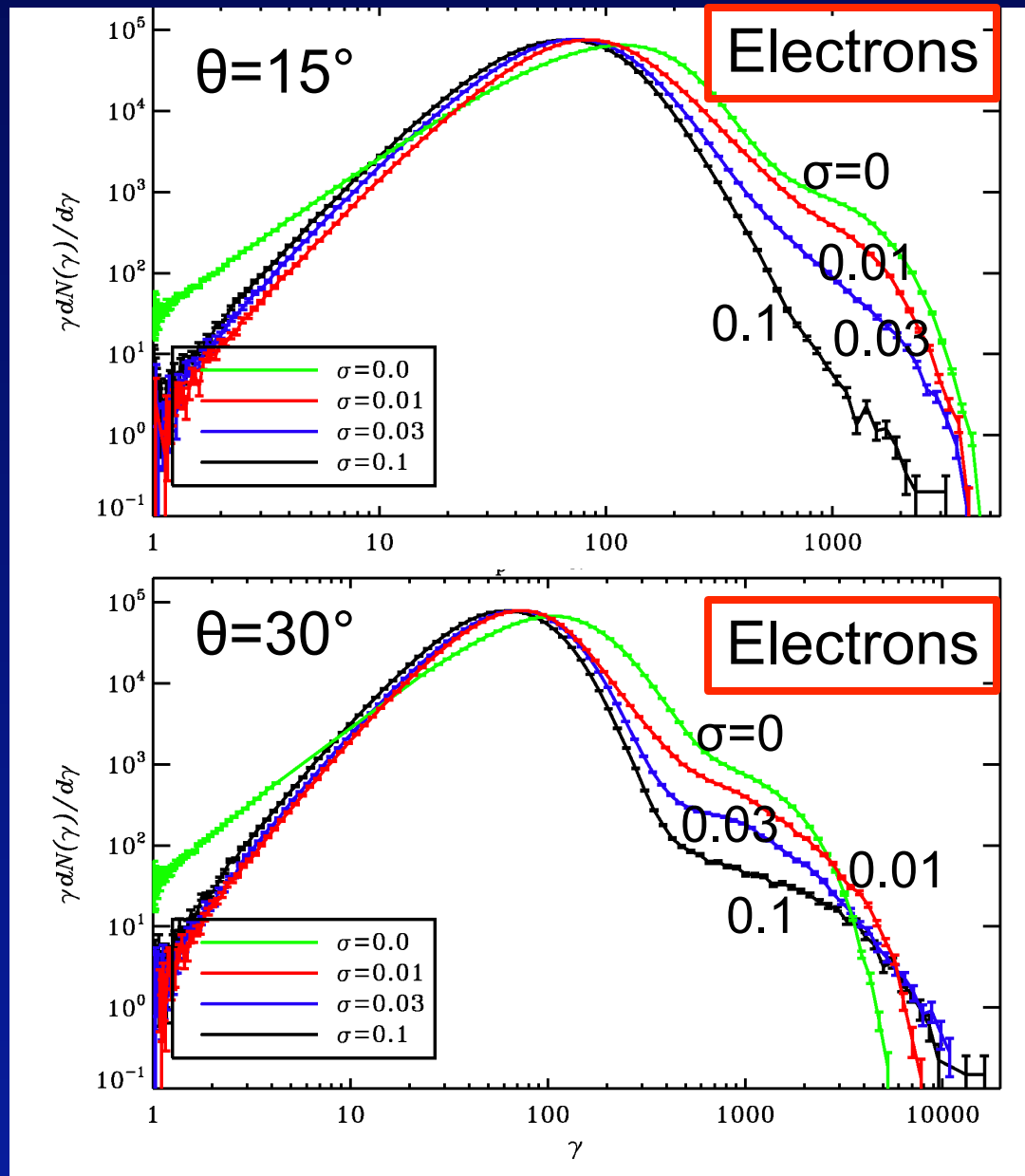


Ion vs electron acceleration (2/2)

$\sigma=0.1$ $\theta=15^\circ$ e⁻-p⁺ shock: **ELECTRONS** are more strongly tied to the magnetic field lines and get quickly advected downstream



Varying σ in e^-p^+ shocks



With **increasing** σ , electrons are **more tied** to magnetic field lines
→ Once advected downstream, it is **harder** for them to come back upstream
→ **Lower** efficiency of electron acceleration, **independent of θ**

Summary

- **Relativistic magnetized** ($\sigma=0.01-0.1$) collisionless shocks do exist
- For fixed magnetic field strength, the shock structure and acceleration efficiency critically depend on the magnetic inclination (**subluminal vs superluminal shocks**):
 - Subluminal **pair** shocks ($\theta < \theta_{\text{crit}} \approx 34^\circ$) are efficient particle accelerators ($\sim 1\%$ by number, $\sim 10\%$ by energy), superluminal shocks are not.
 - **IONS** are efficiently accelerated in **subluminal** electron-ion shocks, with $\sim 3\%$ of particles and $\sim 20\%$ of energy stored in a suprathermal tail.
 - **ELECTRON** acceleration in **subluminal** electron-ion shocks is $\sim 5-10$ times less efficient than for ions, especially for high magnetizations ($\sigma \sim 0.1$).

Implications

Constraints on the composition and magnetization of pulsar winds:

- If **electron-positron** plasma, then nearly-parallel shocks (in the upstream fluid frame $\theta'_{\text{crit}} \approx 34^\circ / \gamma_0$) are required for efficient particle acceleration; or magnetization must be $\sigma \leq 10^{-3}$
- If **electron-ion** plasma, magnetization must be $\sigma \leq 10^{-2}$ regardless of the magnetic obliquity, since for $\sigma \sim 0.1$ shocks, electron acceleration is inefficient **even for subluminal configurations**.

Caveats:

- Long-term shock evolution? Results from 3D simulations? Realistic mass ratios?
- Different magnetic field geometry in the upstream flow: Magnetic turbulence? Striped wind? → Acceleration via reconnection?
- Different composition of the upstream flow: Ion-doped pair plasma? → Acceleration via Resonant Cyclotron Absorption (Elena's talk)?