

Radiative jets in global simulations of super-critical accretion disks in general relativity

Einstein Symposium

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Super-critical accretion

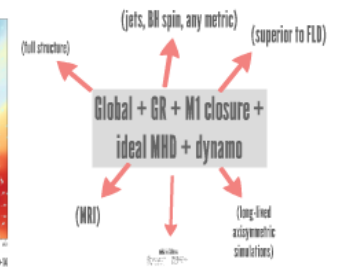
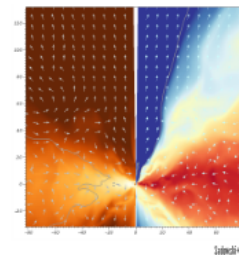


- Spherical accretion limited by the critical (Eddington) luminosity

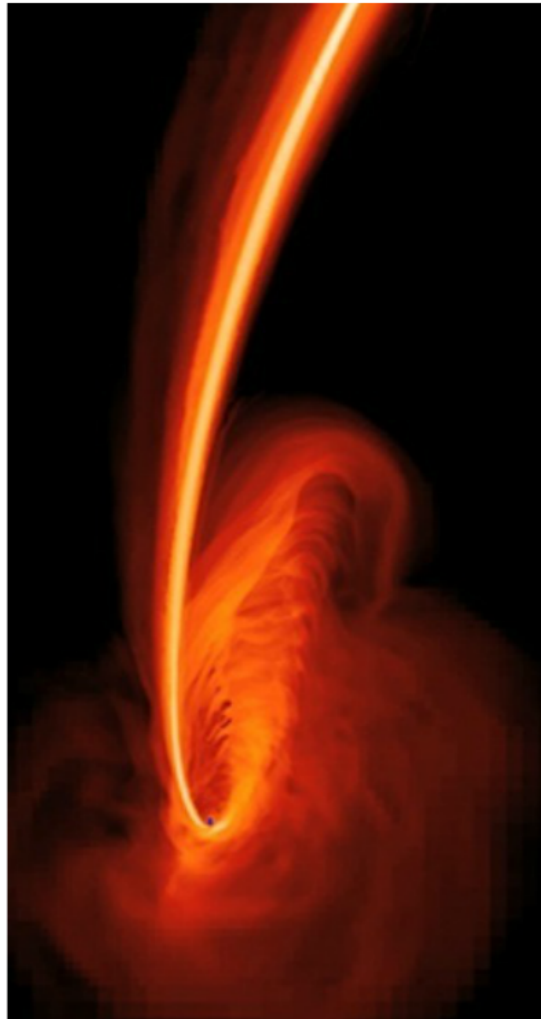
$$L_{\text{Edd}} = \frac{4\pi cGMm_{\text{p}}}{\kappa_{\text{es}}} = 1.25 \times 10^{38} \frac{M}{M_{\odot}} \text{erg/s}$$

- Breaking the spherical symmetry allows for luminosities exceeding the Eddington limit

KORAL code (Sądowski+13,+14)



Super-critical accretion



(c) James Guillochon & Suvi Gezari

- Spherical accretion limited by the critical (Eddington) luminosity

$$L_{\text{Edd}} = \frac{4\pi cGMm_p}{\kappa_{\text{es}}} = 1.25 \times 10^{38} \frac{M}{M_{\odot}} \text{erg/s}$$

- Breaking the spherical symmetry allows for luminosities exceeding the Eddington limit
- For a thin disk (efficiency for spin zero $\eta = 0.057$), this luminosity corresponds to

$$\dot{M}_{\text{Edd}} = \frac{L_{\text{Edd}}}{\eta c^2} = 2.4 \times 10^{18} \frac{M}{M_{\odot}} \text{g/s}$$

- Radiatively inefficient accretion allows for super-critical accretion rates with moderate luminosities
- Whenever there is enough gas close to the BH - super-critical accretion

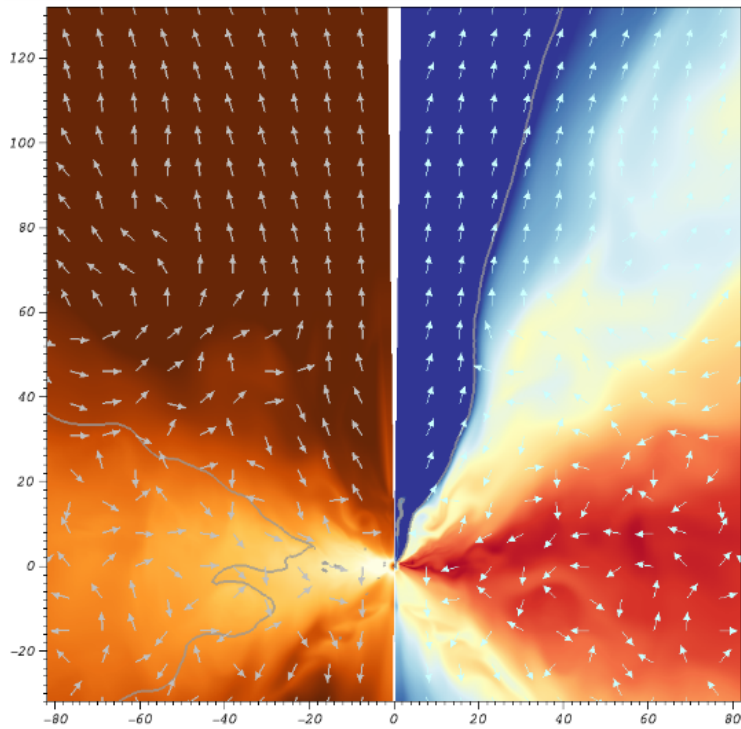
mergers of
galaxies

tidal disruption
events

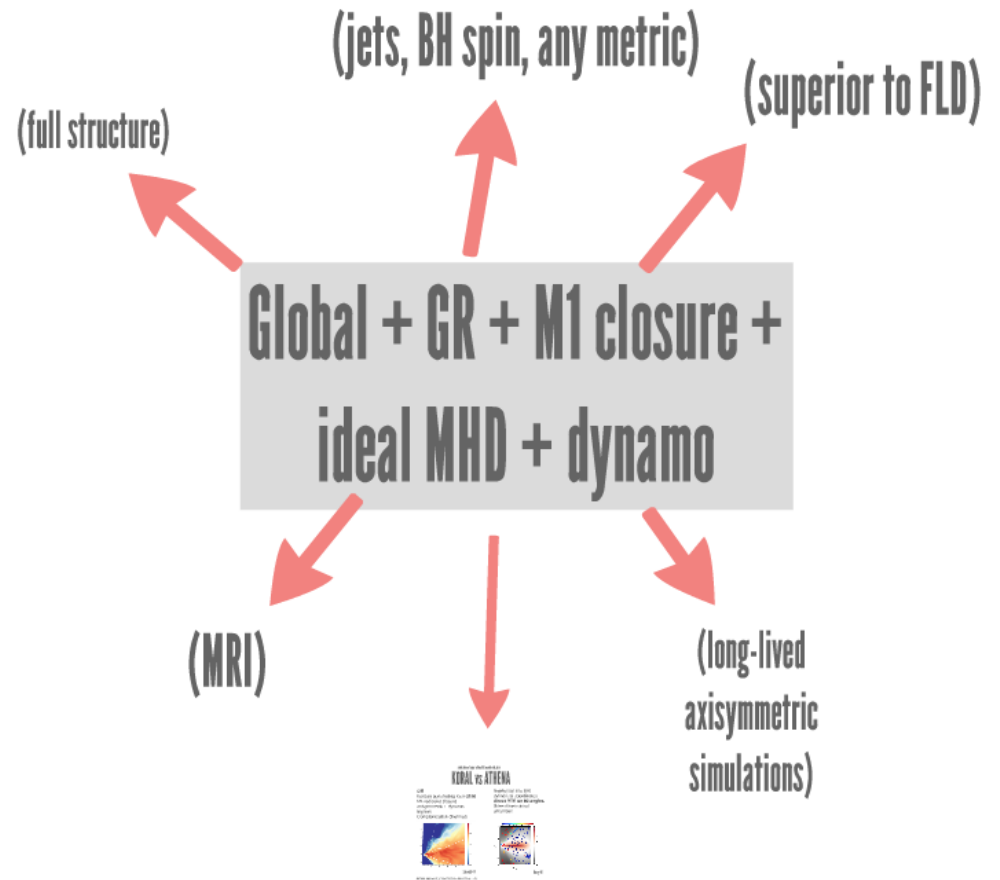
GRBs

ULXs (?)

KORAL code (Sadowski+13,+14)



Sadowski+14



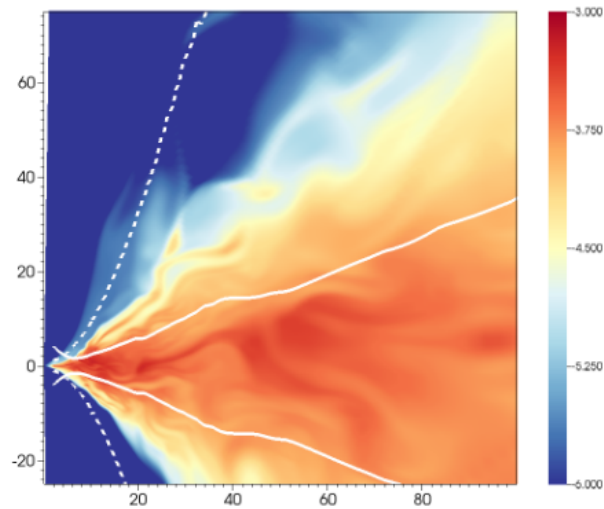
simulations of super-critical BH accretion disks in

KORAL vs ATHENA

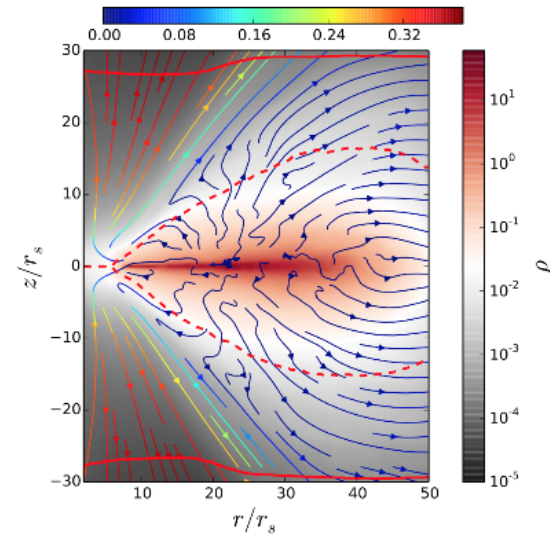
GR

horizon penetrating Kerr-Schild
M1 radiative closure
axisymmetric + dynamo
big box
Comptonization (thermal)

Newtonian (no BH)
cylindrical coordinates
direct RTE on 80 angles
three dimensional
small box



Sadowski+14



Jiang+14

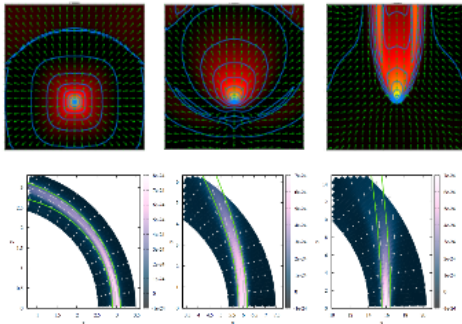
Similar approach in HARMRAD (McKinney+14)
and Cosmos++ (Fragile+14)

M1 a covariant radiative closure

The closure: M1 (En. density & Flux)

There is a frame where flux vanishes and where the radiation stress-energy tensor has only symmetric diagonal terms

$$R_{\text{rf}}^{ij} = \begin{bmatrix} E_{\text{rf}} & 0 & 0 & 0 \\ 0 & \frac{1}{3}E_{\text{rf}} & 0 & 0 \\ 0 & 0 & \frac{1}{3}E_{\text{rf}} & 0 \\ 0 & 0 & 0 & \frac{1}{3}E_{\text{rf}} \end{bmatrix}$$



👍 reasonable, local, simple in GR, cheap
 👎 far from being perfect

Radiation - gas coupling

Conservation of mass, energy & momentum:

$$\begin{aligned} (\rho u^\mu)_{;\mu} &= 0 \\ (T_\nu^\mu)_{;\mu} &= G_\nu \\ (R_\nu^\mu)_{;\mu} &= -G_\nu \end{aligned}$$

Radiative four-force:

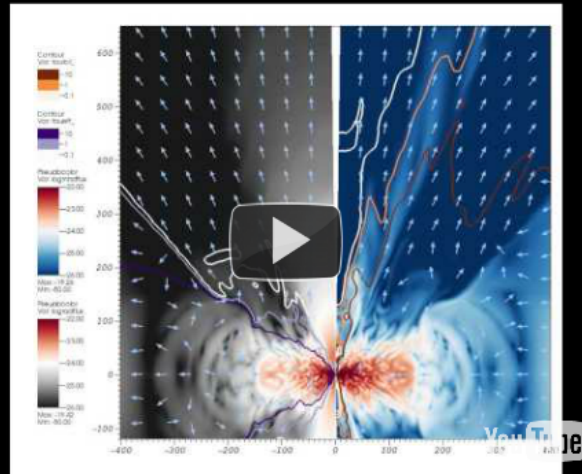
$$\hat{G}^\mu = \begin{bmatrix} \rho \kappa_{\text{abs}} (\hat{E} - 4\pi B) \\ \rho (\kappa_{\text{abs}} + \kappa_{\text{es}}) \hat{F}^i \end{bmatrix}$$

Thermal Comptonization:

$$\hat{G}_{\text{Compt}}^t = \rho \kappa_{\text{es}} \hat{E} \frac{4k}{m_e} (\hat{T}_{\text{rad}} - \hat{T}_{\text{gas}}) \left(1 + \frac{4k}{m_e} \hat{T}_{\text{gas}} \right)$$

Super-critical disk

spin zero, accretion rate 30 times Eddington



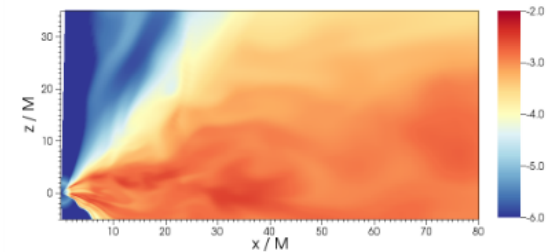
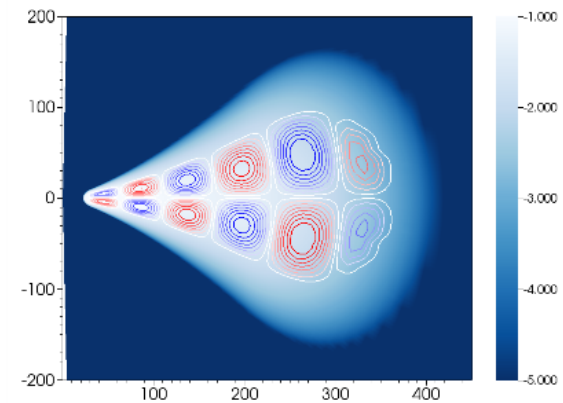
jets in super-critical disks simulation setup

- initiated as equilibrium torus of gas and radiation
- supermassive black hole
- multiple loops of initial magnetic field
 - weak magnetic flux limit
- zero BH spin
 - no Blandford - Znajek
- wide range of super-Eddington accretion rates

Table 1. Model parameters

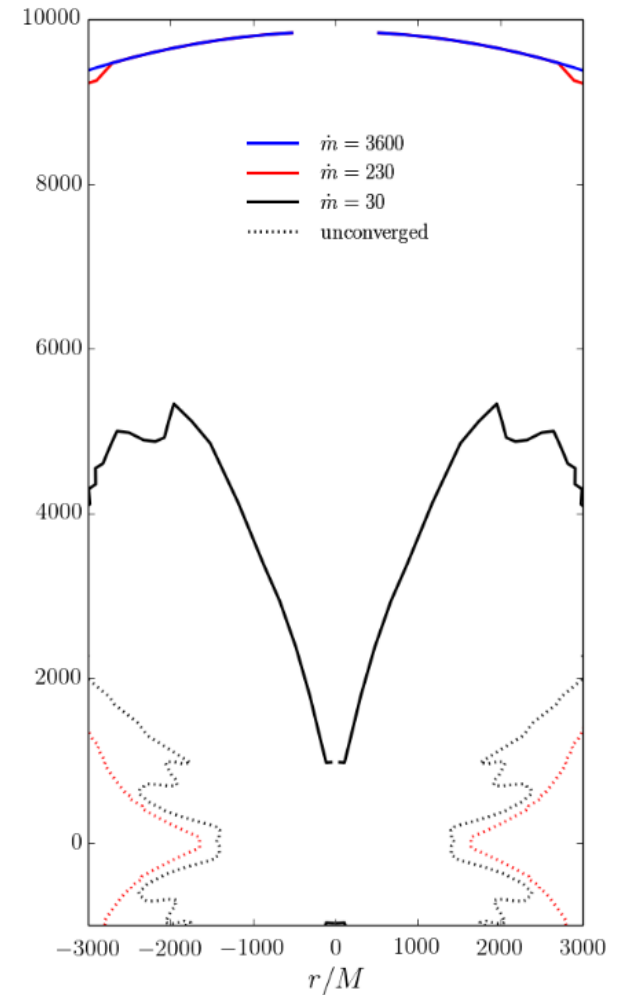
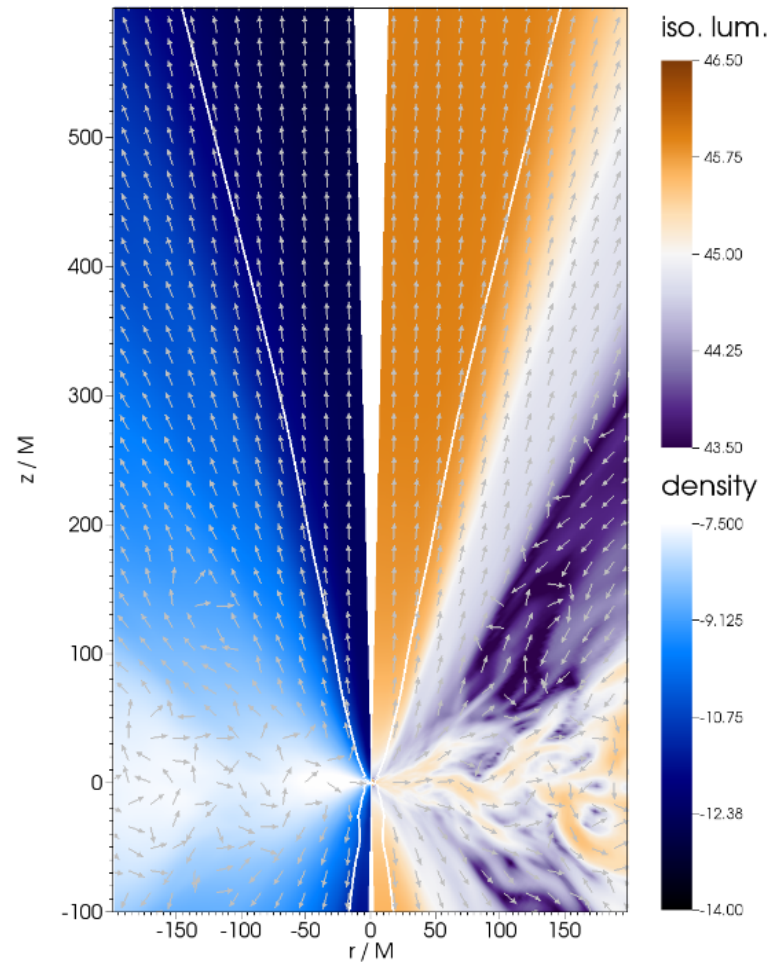
Name	\mathcal{K}	$t_{\max}/(GM/c^3)$	$\langle \dot{M} \rangle / \dot{M}_{\text{Edd}}$
A	10.0	78,000	30
B	5.0	91,000	230
C	1.0	140,000	3600

Other parameters: $M_{\text{BH}} = 3 \times 10^5 M_{\odot}$, $a_* = 0.0$, resolution: 304x192, $R_{\min} = 1.85$, $R_{\max} = 10000$, $R_0 = 1.0$, $H_0 = 0.6$, $\beta_{\max} = 10.0$. All definitions from Sądowski et al. (2014b).

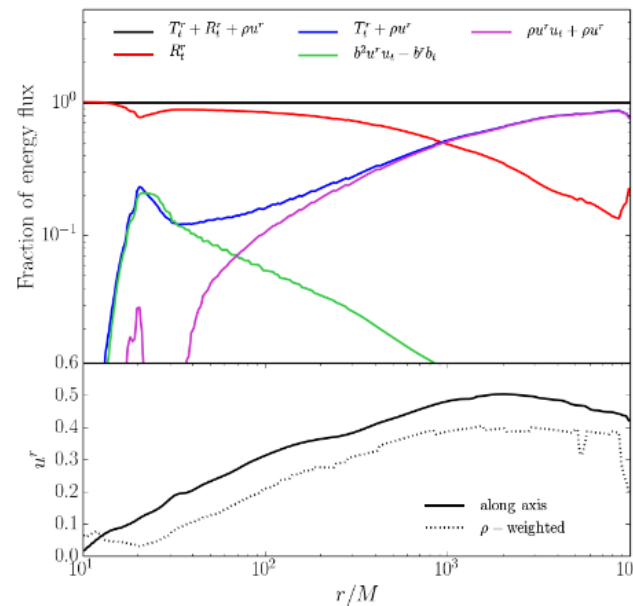
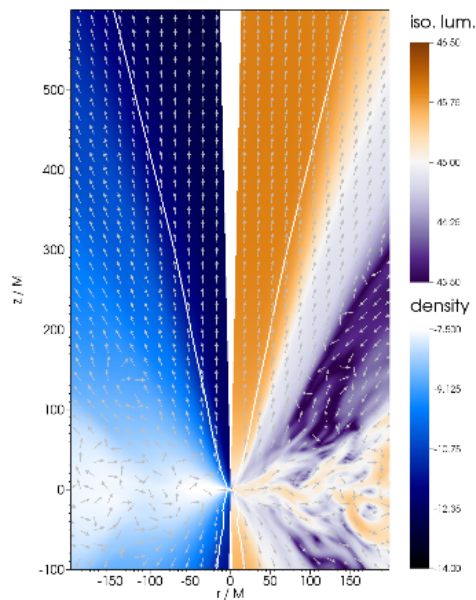


Funnel

- accretion of gas in a thick disk
- outflows of mass at $R > 20R_g$
- low density funnel region
- but the photosphere far from the BH
- energy flows out along the axis
- super-Eddington flux

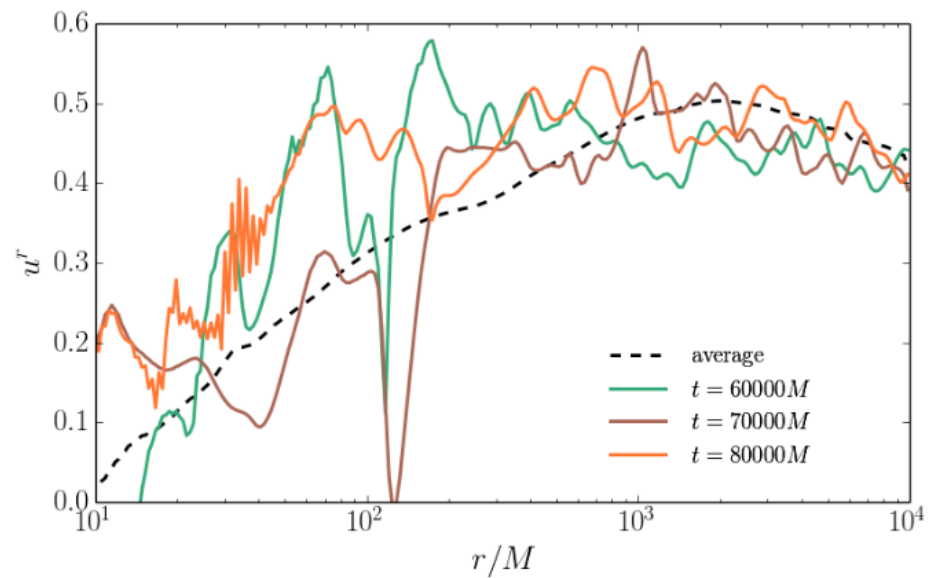
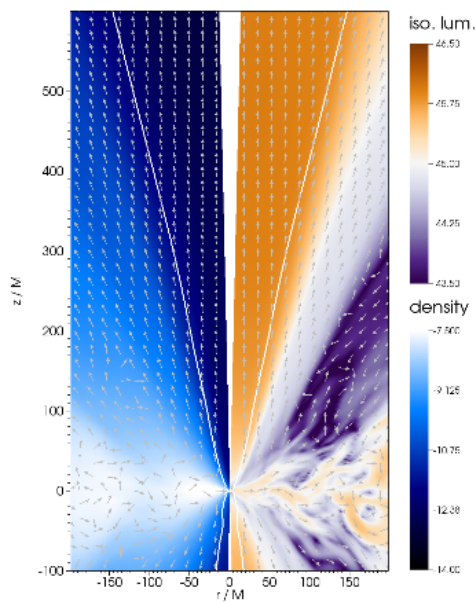


Energy fluxes in the funnel



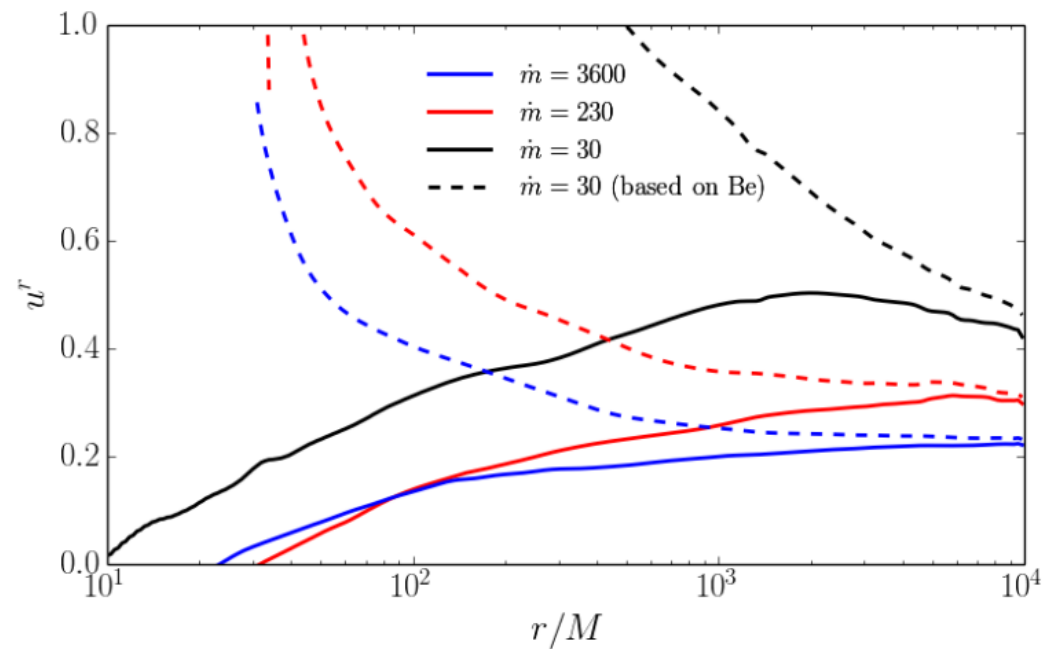
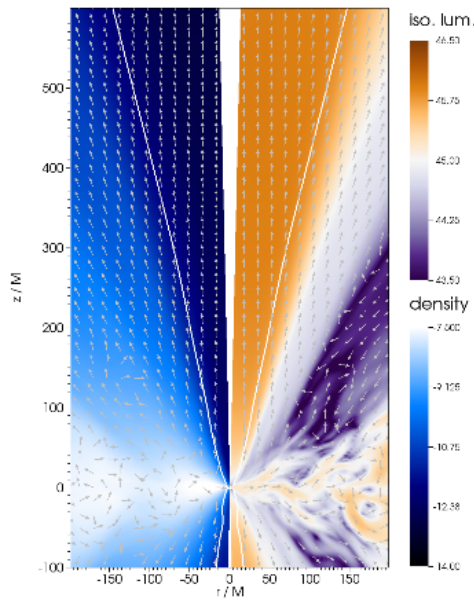
- Luminosity in the funnel
~1% of accreted rest mass energy
- Gas unbound
- Energy flux dominated initially by the radiative flux, converting into kinetic energy of gas as long as it is optically thick
- Gas reaches mildly-relativistic velocities ~0.5c

Gas velocity in the jet



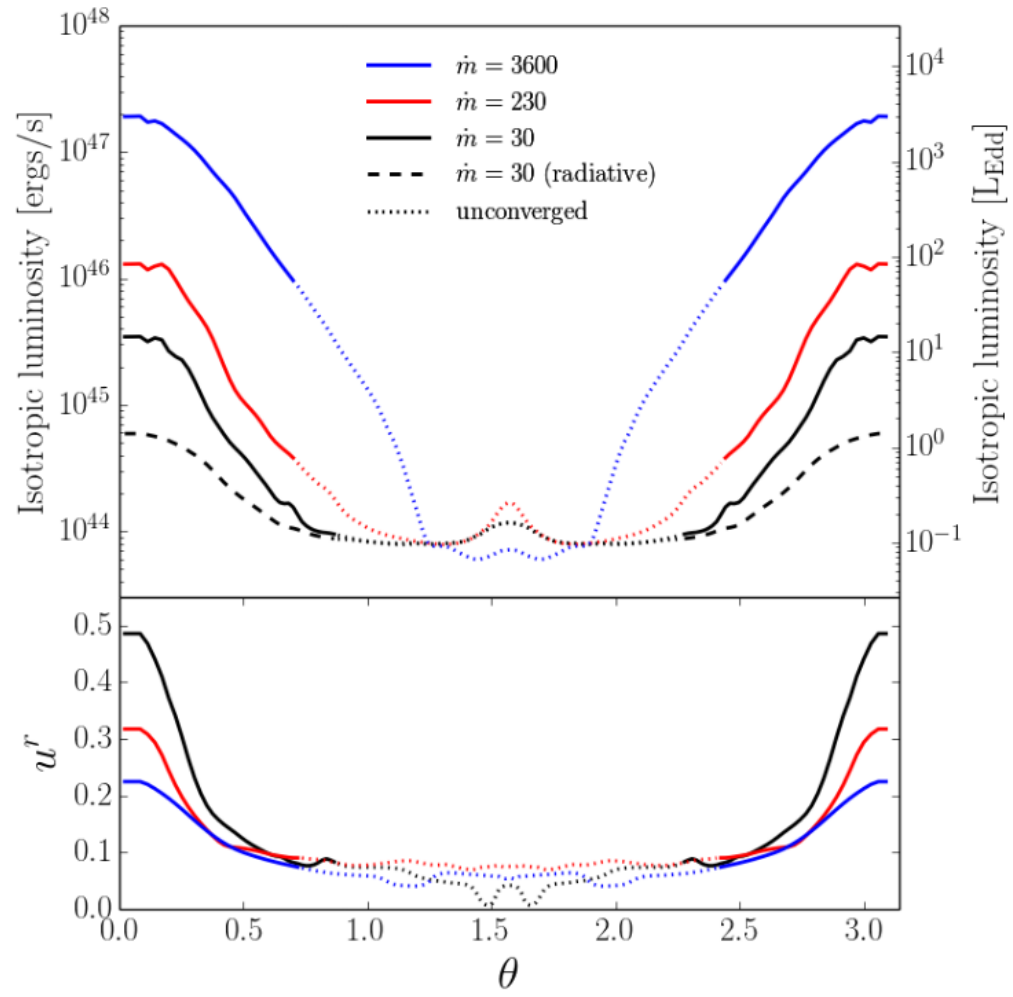
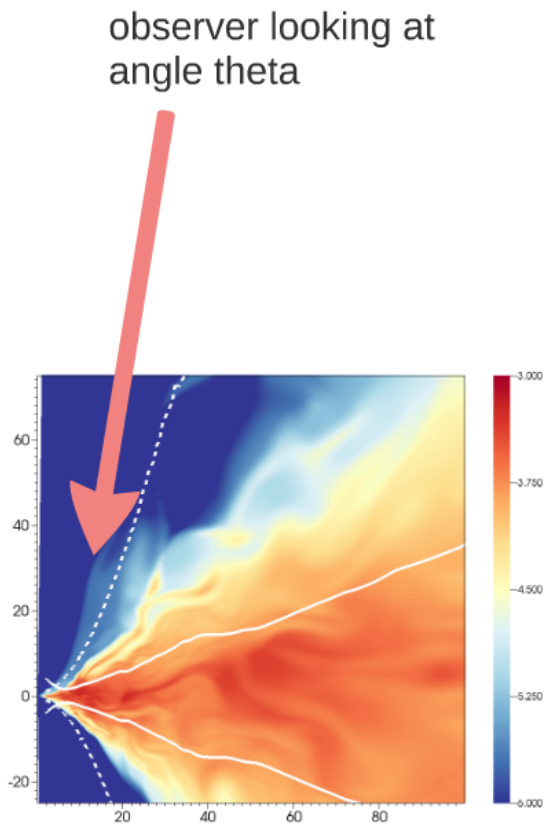
- Gas velocities well resolved in the funnel
- Non-uniform
- Fluctuations may lead to shocks once the gas is optically thin

Gas velocity in the jet



- Higher accretion rates imply higher gas densities
- Larger optical depth in the funnel
- Lower gas velocities

Isotropic equivalent luminosities



Tidal Disruption Events

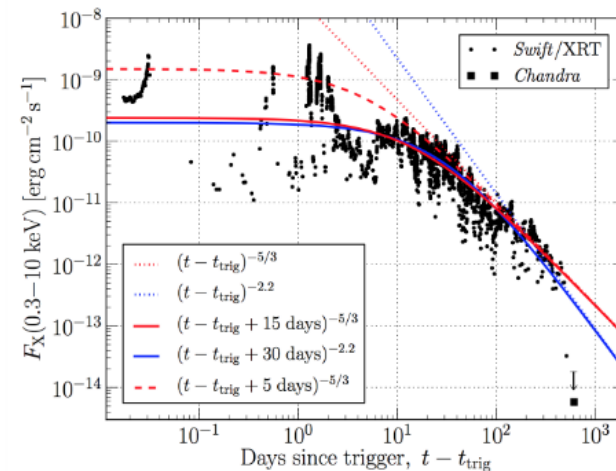


(c) James Guillochon & Suvi Gezari

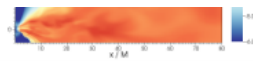
- whenever a star gets too close to a SMBH it gets disrupted
- accretion coming from the fallback of the bound material
- accretion rate starting at ~ 1000 Eddington
- followed by a $t^{-5/3}$ decay

Swift J1644+57

- isotropic equivalent luminosity in X-ray $\sim 1e48$ erg/s
- in the standard picture jet emission requires significant accumulation of magnetic flux and non-zero BH spin
- radio afterglow suggests Lorentz factor $\sim 2-5$



Tchekhovskoy+14



Radiative jets in GR simulations of super-critical BH accretion disks

- Super-critical accretion produces powerful jet-like outflows even without BZ!
- Isotropic equivalent luminosity up to $L_{\text{iso}} \lesssim 10^{48}$ erg/s for $10^6 M_{\odot}$ and $1000 \dot{M}_{\text{Edd}}$
- No magnetic flux or BH spin required
- Outflow radiatively driven
- Collimated by optically thick outflow emerging from the innermost region
- Provides enough energy but too small Lorentz factor for TDEs

